

TNS and non-equilibrium dynamics

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Miguel Frías-Pérez, Luca Tagliacozzo (IFF)

PRB 106, 115117 (2022)

arXiv:2308.04291



MAX PLANCK INSTITUTE
OF QUANTUM OPTICS

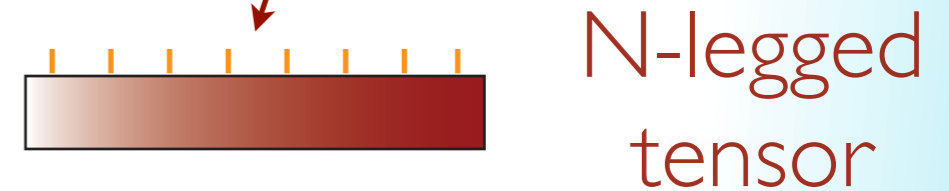


TNRG 31.10.2023

TNS in quantum many-body systems

A general state of the N -body Hilbert space has exponentially many coefficients

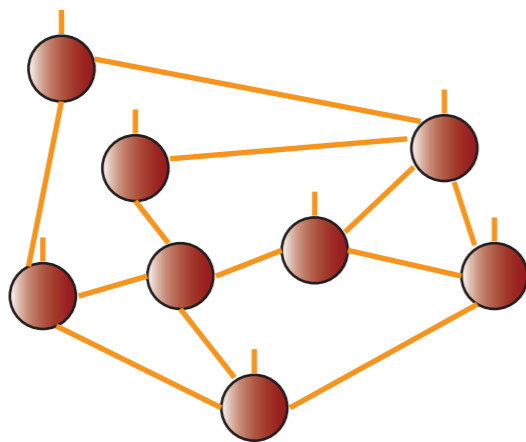
$$|\Psi\rangle = \sum_{i_j} c_{i_1 \dots i_N} |i_1 \dots i_N\rangle$$



A TNS has only a polynomial number of parameters

$$d^N$$

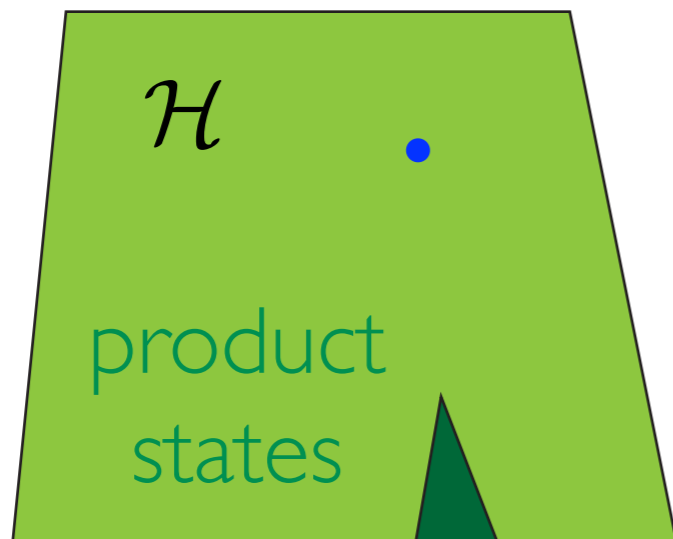
$\text{poly}(N)$



WHY SHOULD TNS BE USEFUL?

States appearing in Nature are peculiar

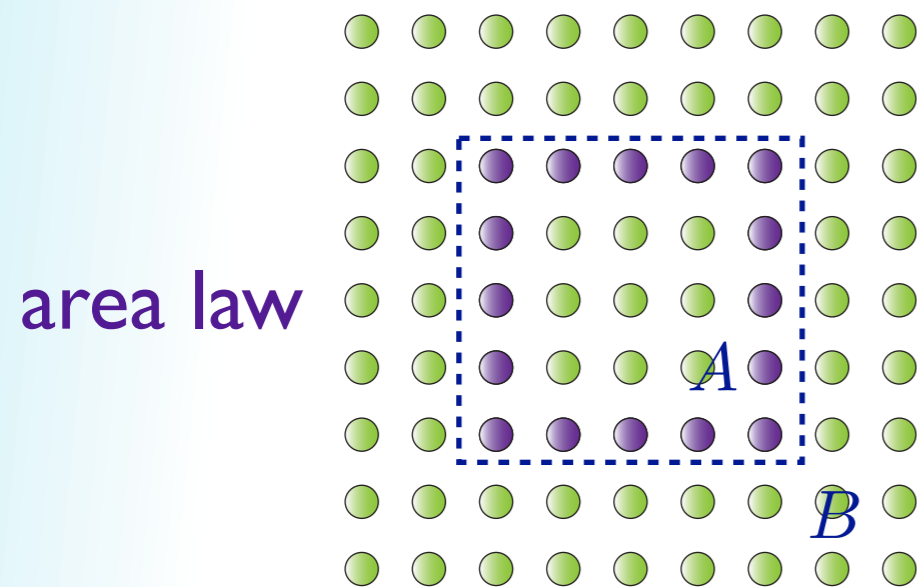
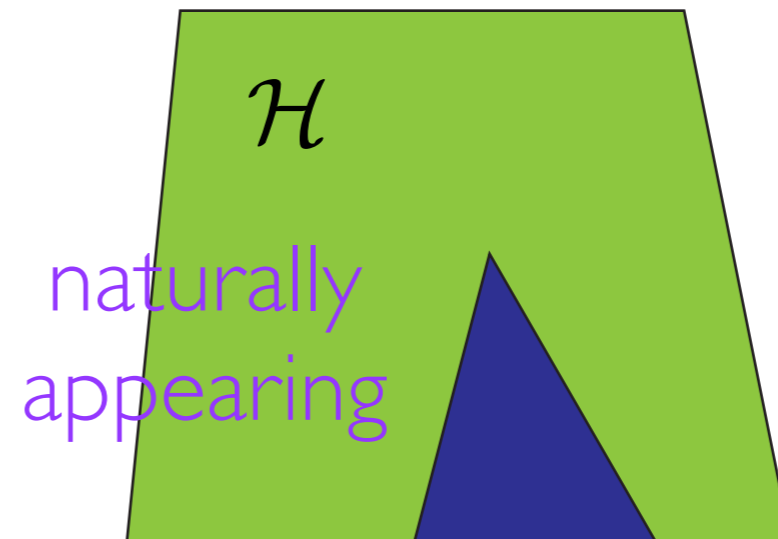
State at random from
Hilbert space is not
close to product



WHY SHOULD TNS BE USEFUL?

States appearing in Nature are peculiar

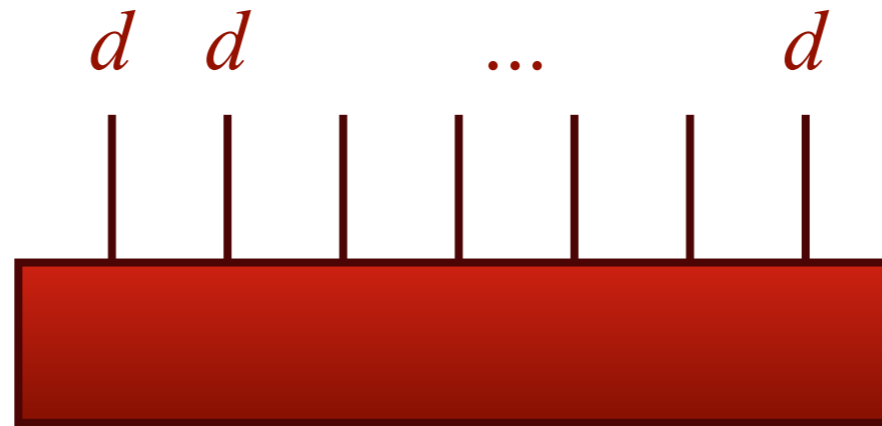
State at random from Hilbert space is not close to product



We look for states with *little* entanglement

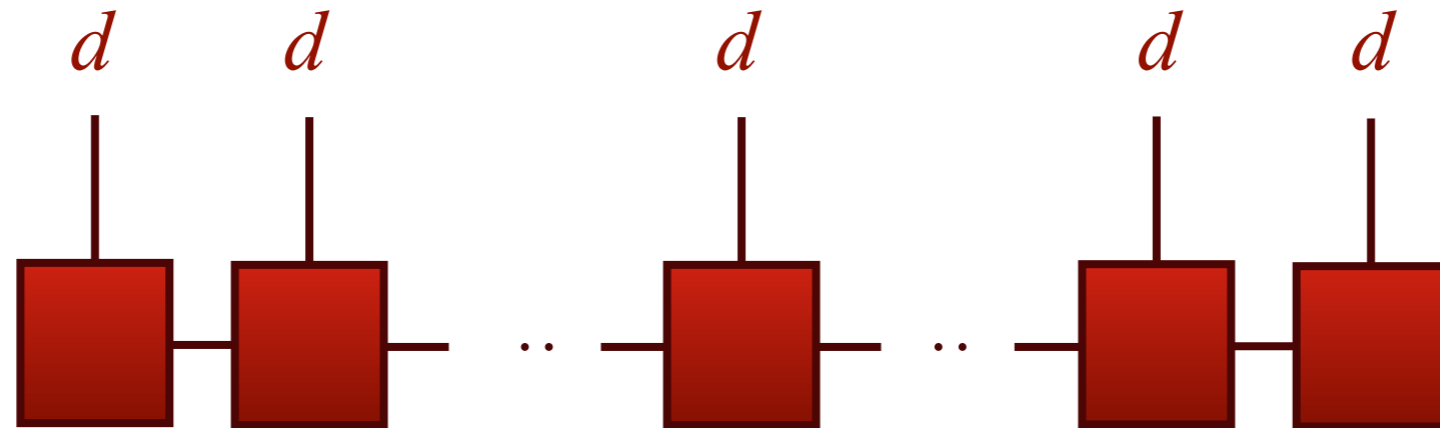
TNS = entanglement based ansatz

MPS

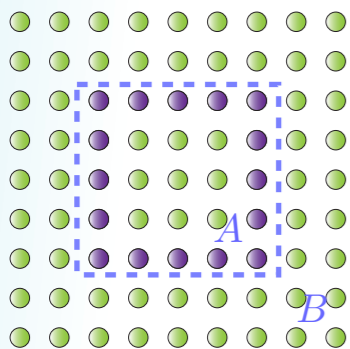


$$|\Phi\rangle = \sum_{s_1, \dots, s_N=1}^d c_{s_1, \dots, s_N} |s_1, \dots, s_N\rangle$$

MPS



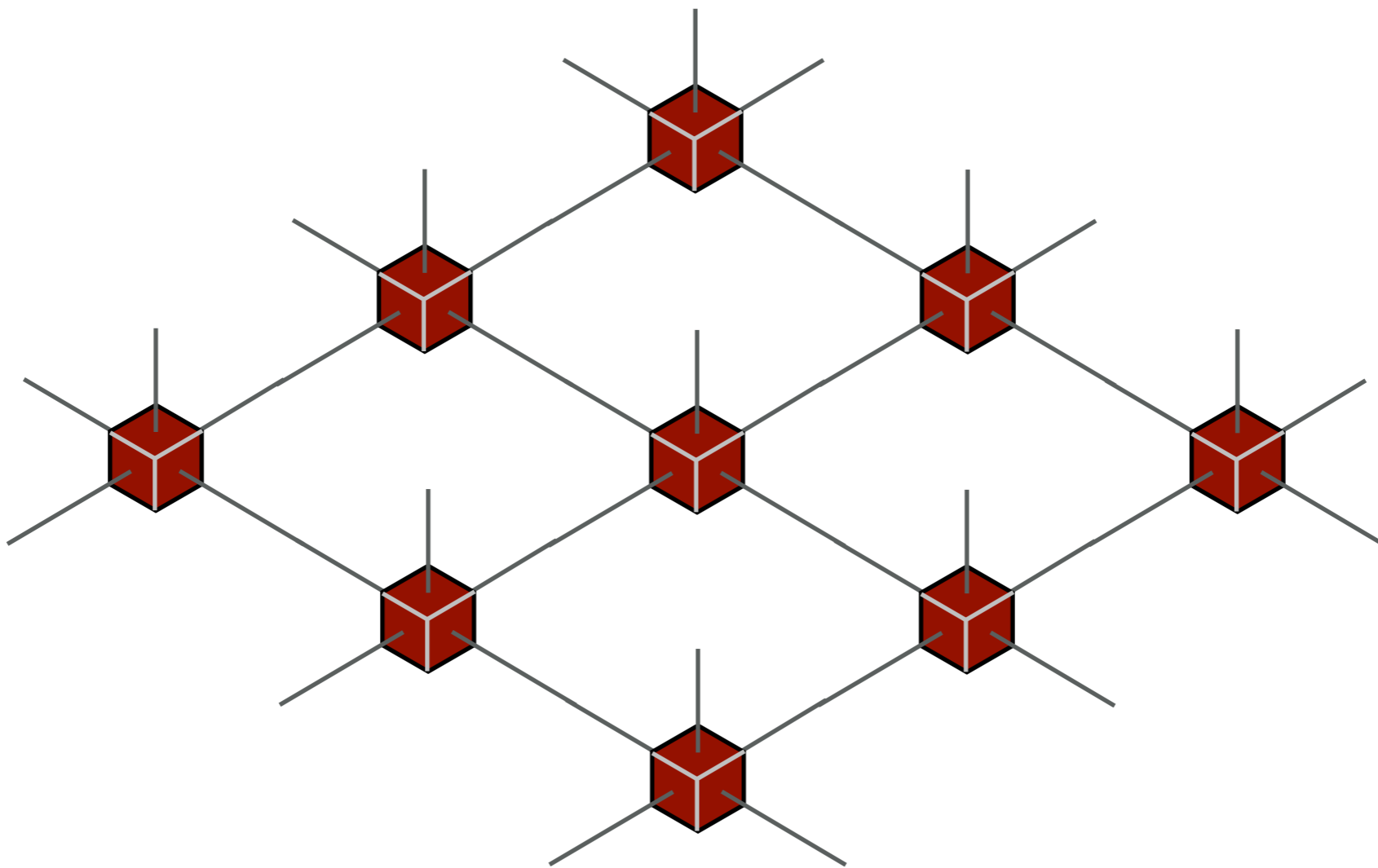
$$|\Phi_D\rangle = \sum_{s_1, \dots, s_N=1}^d \text{tr} (A^{s_1} [1] \dots A^{s_1} [N]) |s_1, \dots, s_N\rangle$$



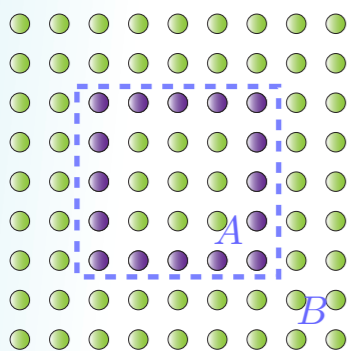
area law by construction

$$S(L/2) \leq \log D$$

PEPS

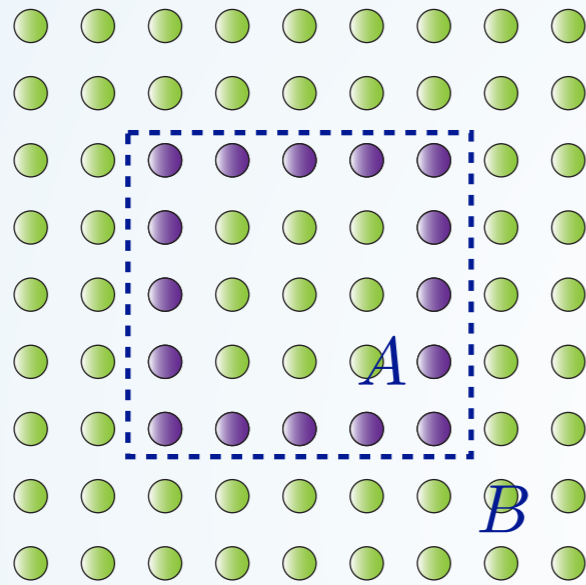


area law by construction

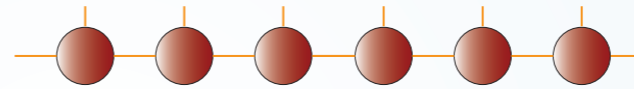


TNS = entanglement based ansatz

area law

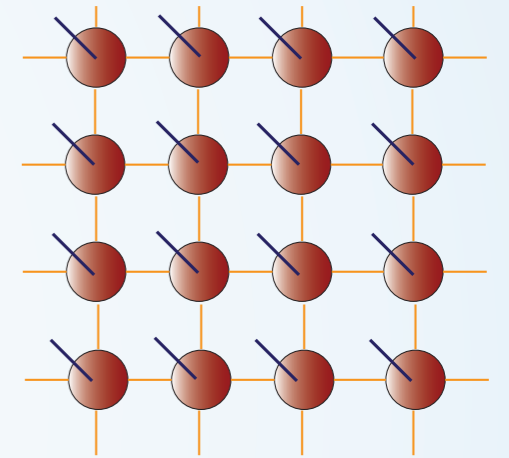


MPS



Schollwöck Ann.Phys.2011

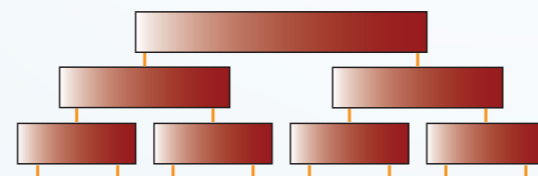
PEPS



Verstraete et al. Adv. Phys. 2008

other TNS

TTN

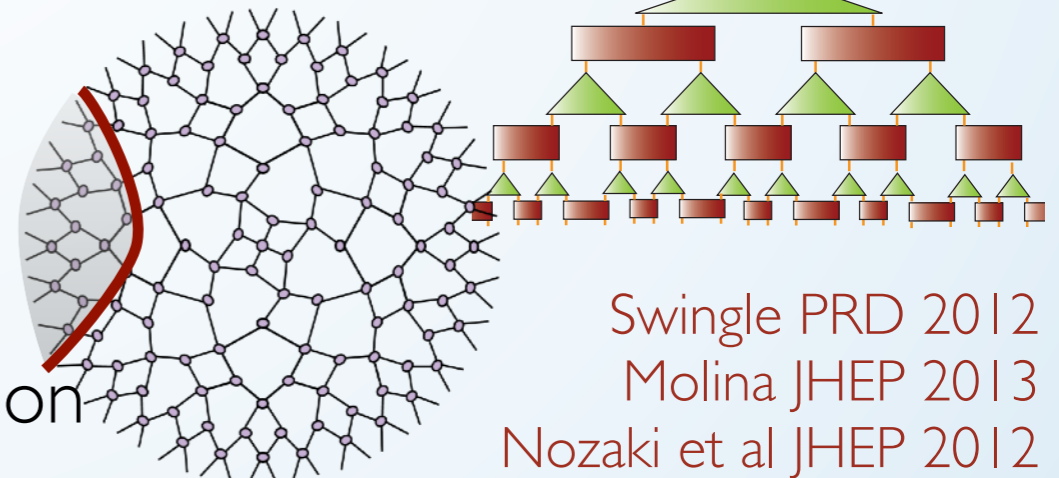


Shi et al PRA 2006

suggested connection
to AdS/CFT

Vidal PRL 2007

MERA



Swingle PRD 2012
Molina JHEP 2013
Nozaki et al JHEP 2012
Bao et al PRD 2015

TNS are very useful in the
quantum many-body context...

works for GS, low energy, thermal equilibrium...

Verstraete, Cirac, PRB 2006 Hastings PRB 2006
Hastings J. Stat. Phys 2007 Molnar *et al.* PRB 2015

area laws

but not for high energy eigenstates, quenches...

Osborne, PRL 2006 Vidmar *et al.*, PRL 2017
Schuch *et al.*, NJP 2008

volume law

entanglement growth in non-equilibrium
scenarios limits the applicability of MPS

fundamental questions: thermalization, ETH...

global quench
in 1D

entanglement
barrier

TNS challenge:
getting around this
limitation

$$D_{\min}(t) \sim e^{\alpha t}$$

Osborne, PRL 2006
Schuch et al., NJP 2008

$$S(t) \propto t$$

some recent progress

Dubail JPhysA 2017
Leviatan et al. 2017
White et al PRB 2018
Surace et al. 2018
Kvornig et al 2021
Rakovzsky et al 2022

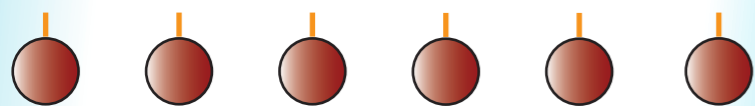


tools to get
dynamical
properties

$t = 0$

$t = \infty$

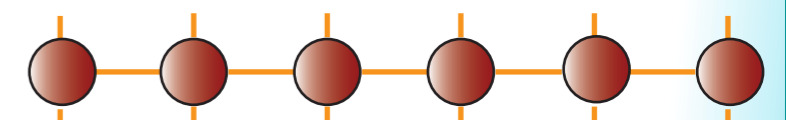
product state



easy to write as MPS

local
observables

thermal states



well approximated as MPO

alternative: give up description of the full state

① light-cone TN for
non-equilibrium
evolution of local
observables

**M. Frías-Pérez, MCB,
PRB 106, 115117 (2022)**

② transforming long-range
entanglement into mixture

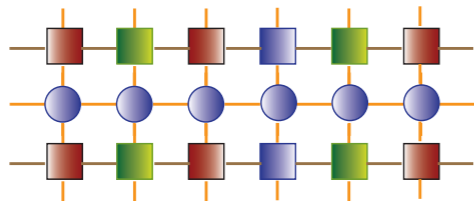
**M. Frías-Pérez, L. Tagliacozzo, MCB,
arXiv:2308.04291**

③ spectral properties of
the QMB Hamiltonian

Yang et al. PRL 124, 100602 (2020),
Lu, MCB, Cirac, PRX Quantum 2, 02032 (2021)
Yang, Cirac, MCB, PRB 106, 024307 (2022)

give up description of the
full state: (local) operators

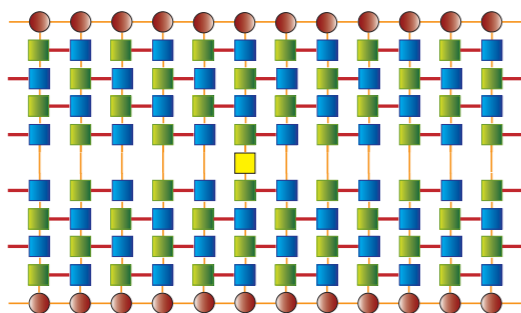
evolving operators: Heisenberg picture Hartmann et al, PRL 2009



also for mixed states
operator space entanglement

Prosen Pizorn, PRL 2008

observables as TN to contract



different *entanglement* quantities

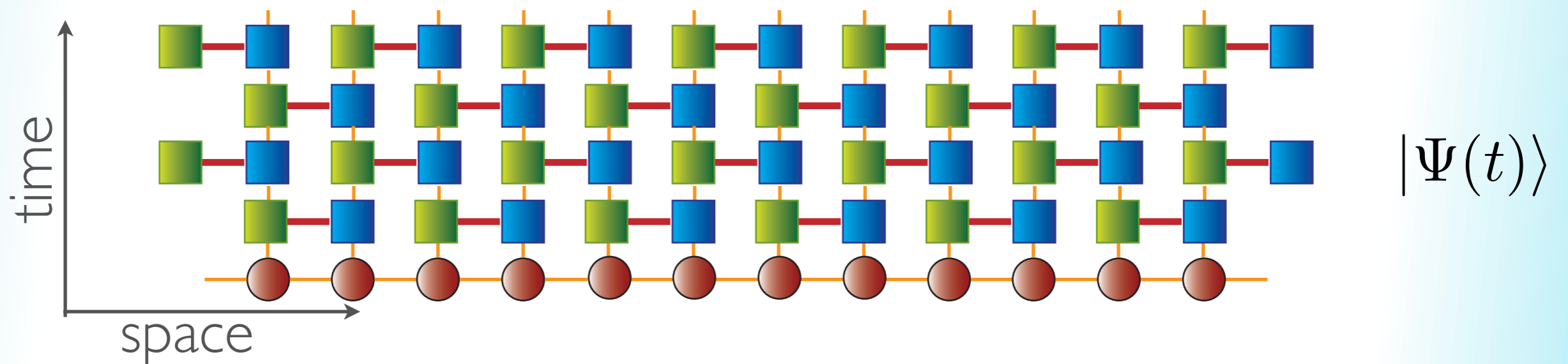
MCB, Hastings, Verstraete, Cirac, PRL 2009

Müller-Hermes et al., NJP 2012

Hastings, Mahajan 2014

Frías-Pérez, MCB PRB 2022

time-dependent observable as a TN

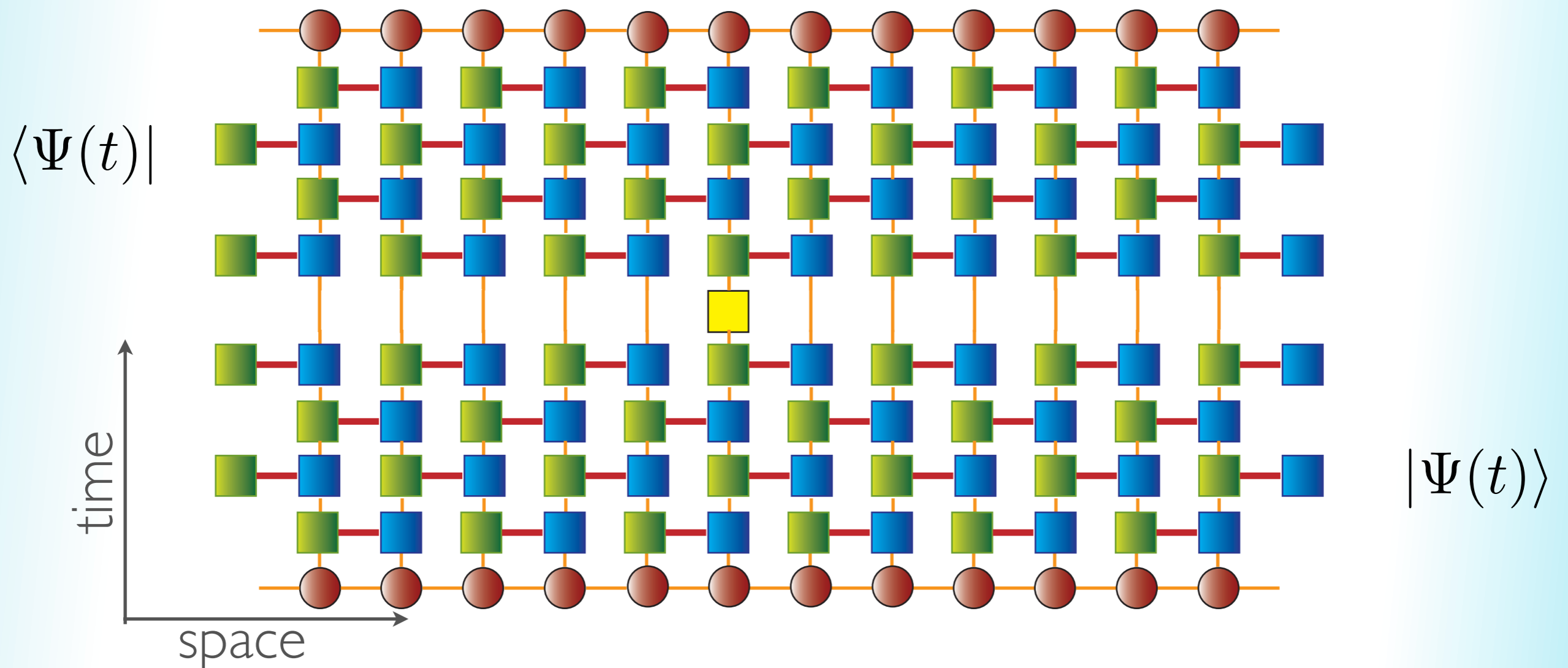


time-dependent observable as a TN

TN describe
observables, not
states

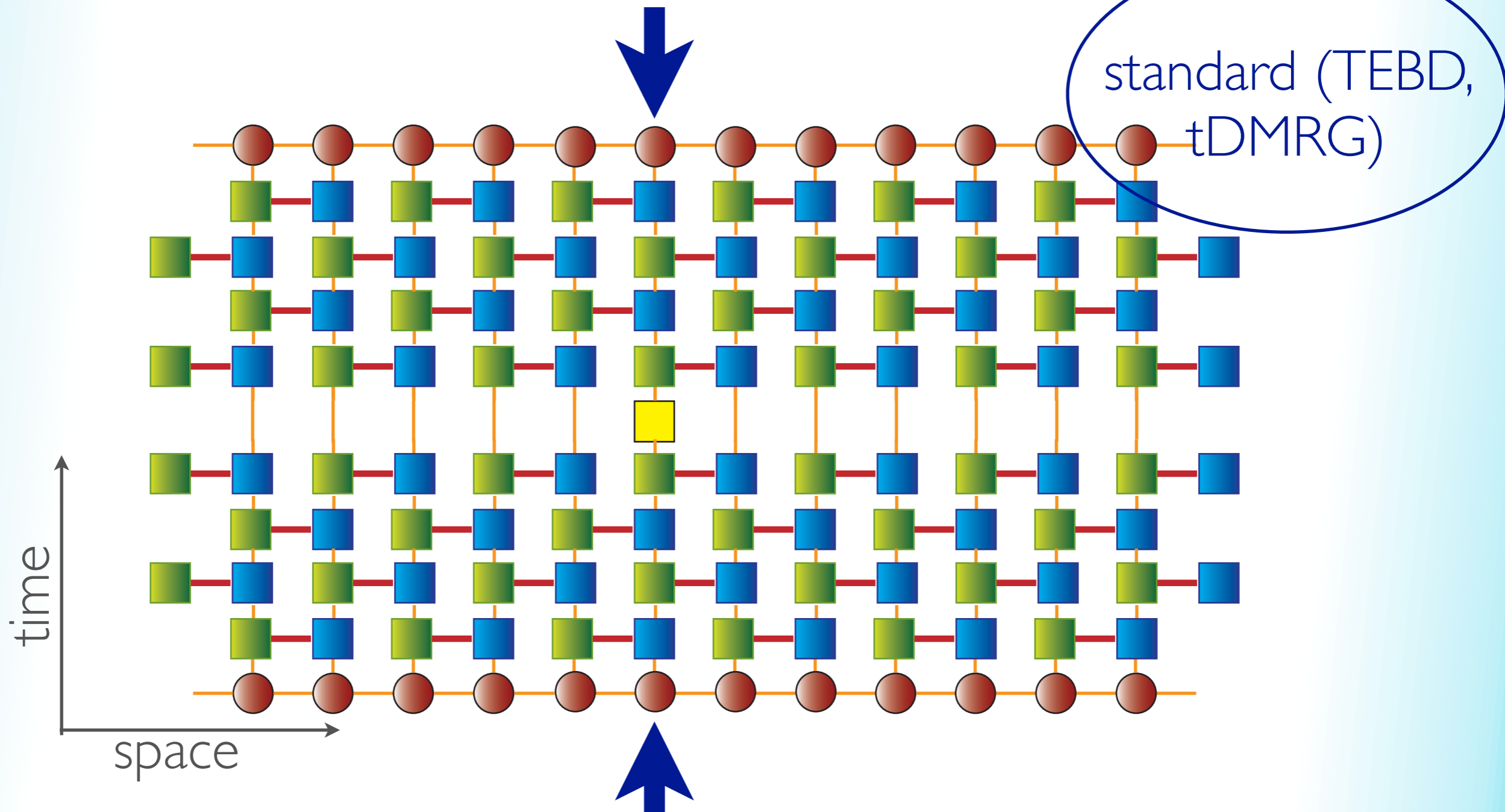
$$\langle \Psi(t) | O | \Psi(t) \rangle$$

exact contraction
not possible
#P complete



time-dependent observable as a TN

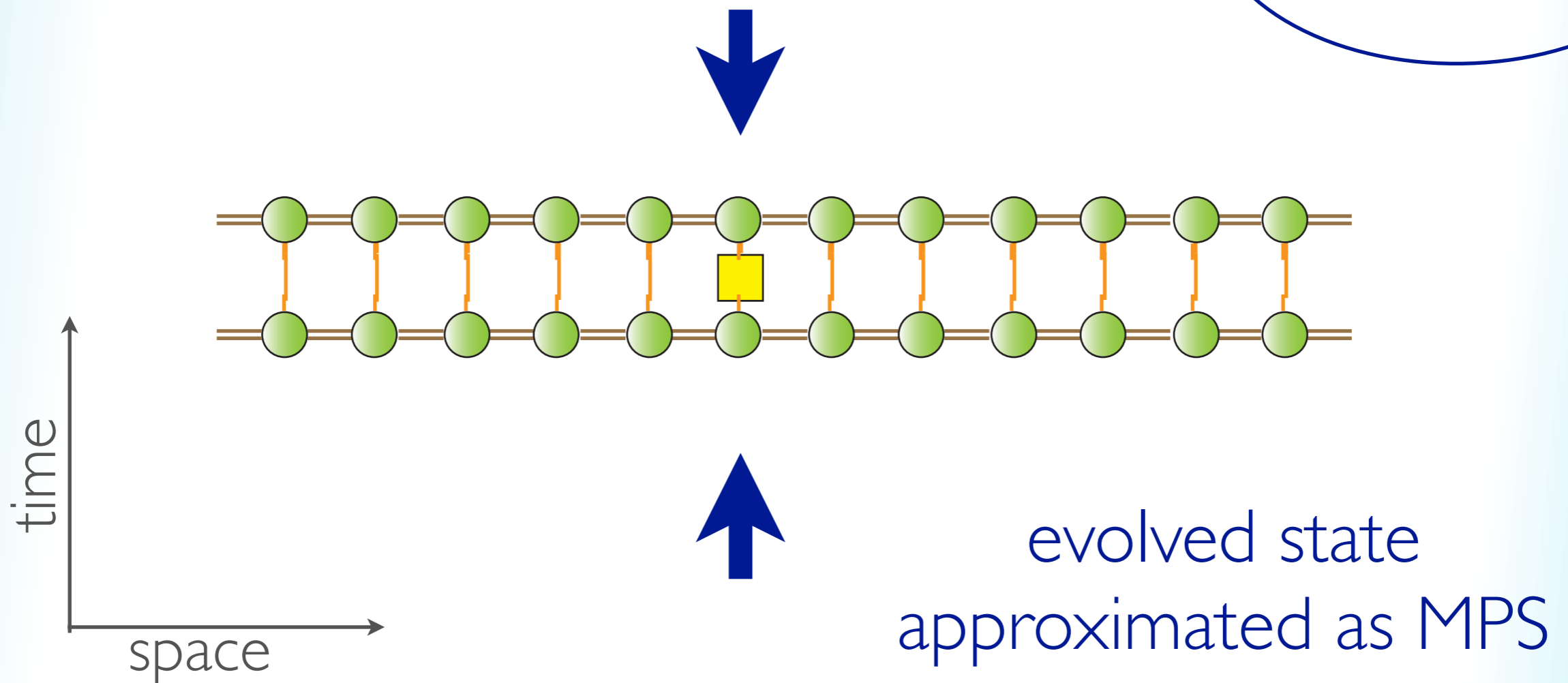
different approximate contraction strategies



time-dependent observable as a TN

different approximate contraction strategies

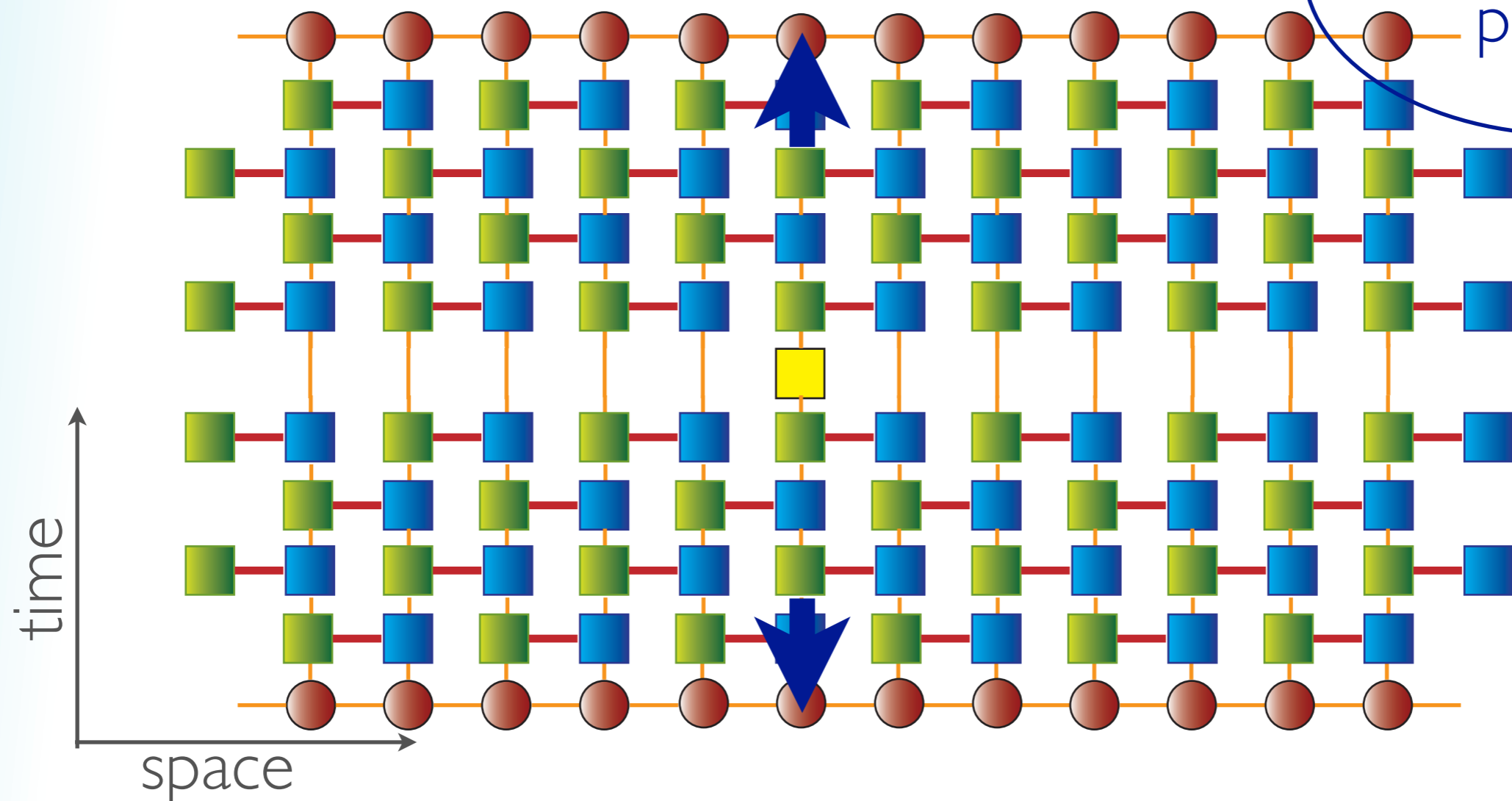
standard (TEBD, tDMRG)



time-dependent observable as a TN

different approximate contraction strategies

evolved operator as
MPO



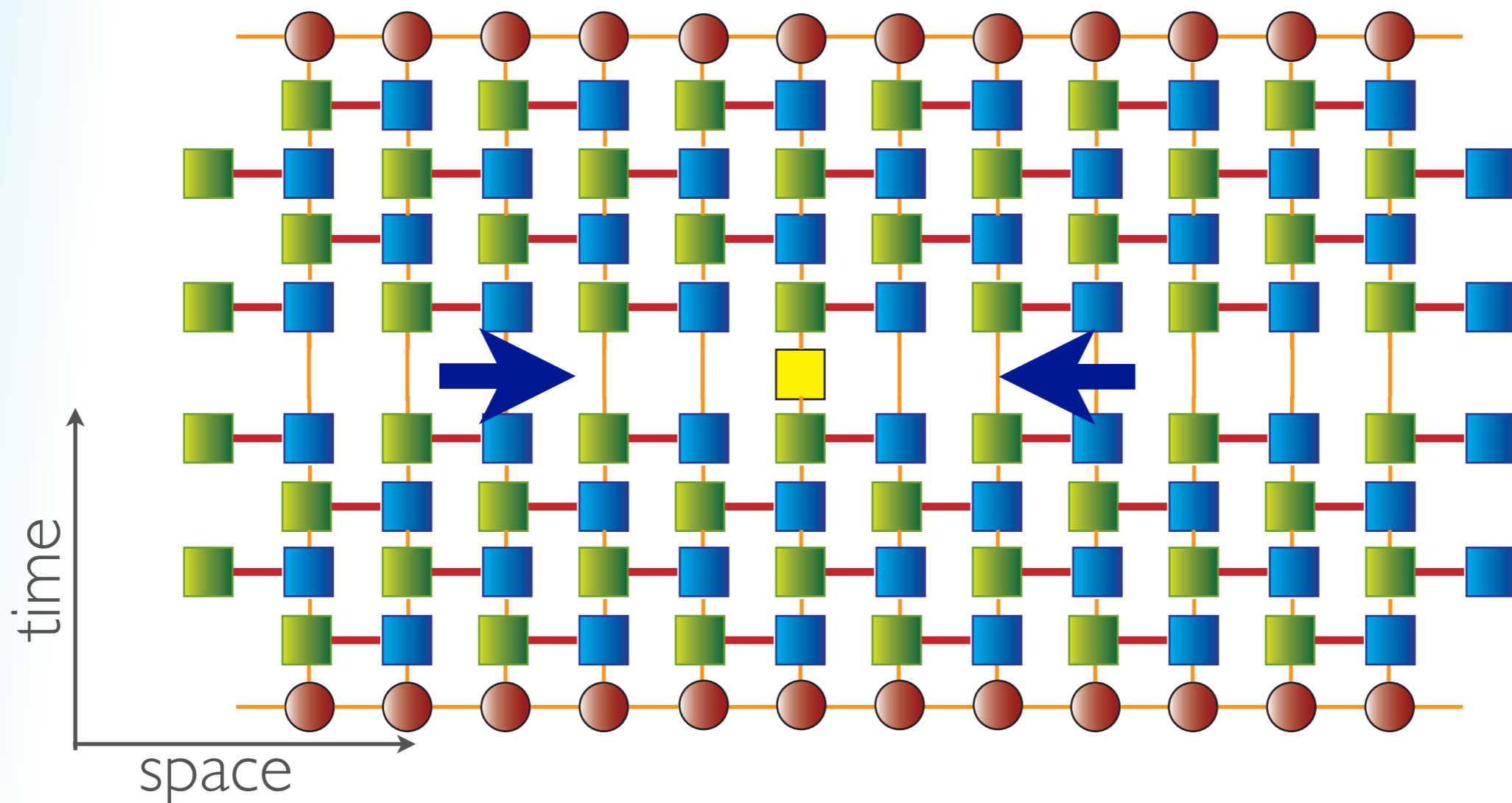
time-dependent observable as a TN

for infinite systems, transverse folding approach

MCB, Hastings, Verstraete, Cirac, PRL 2009

Müller-Hermes, Cirac, MCB, NJP 2012

Hastings, Mahajan 2014



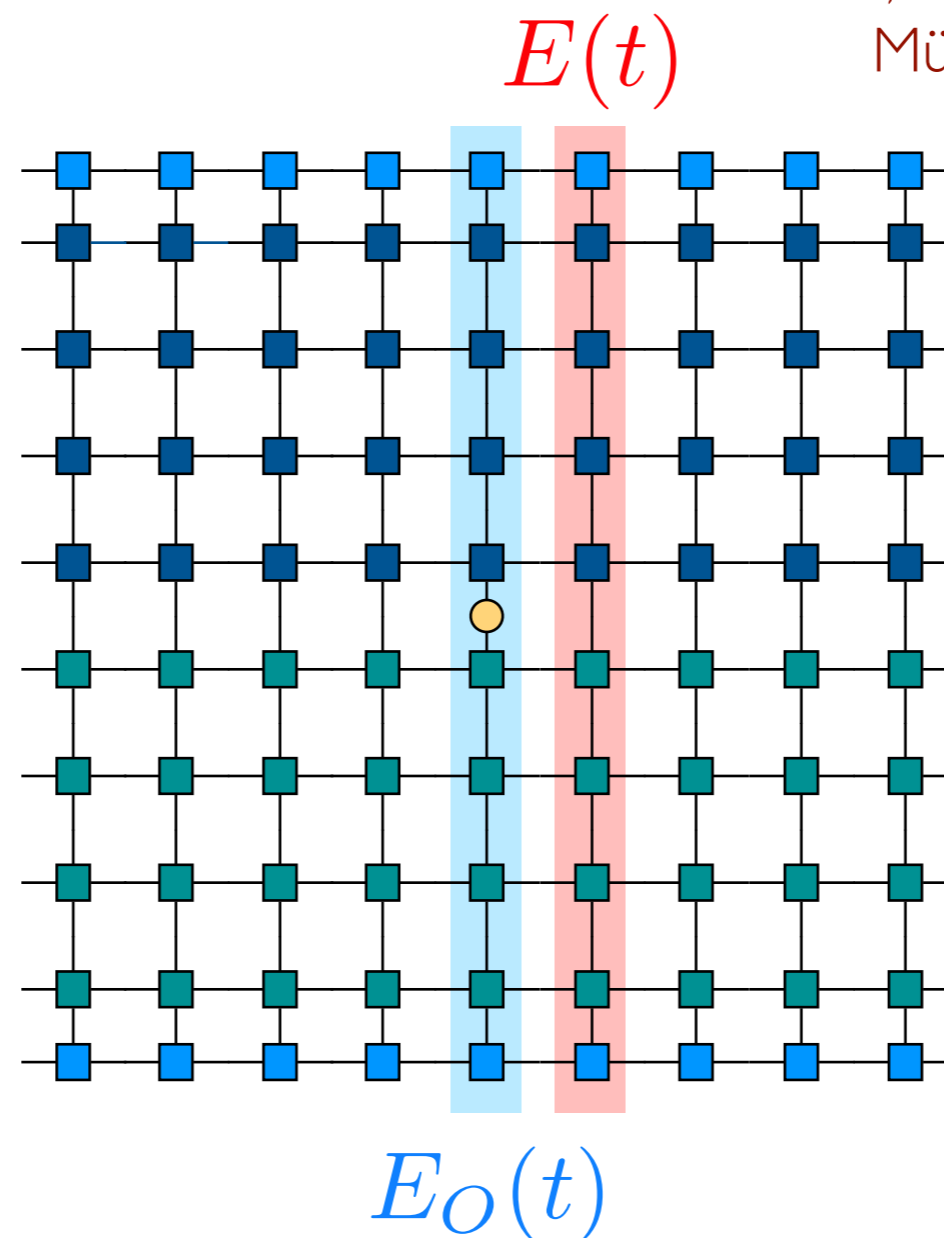
transverse folding approach

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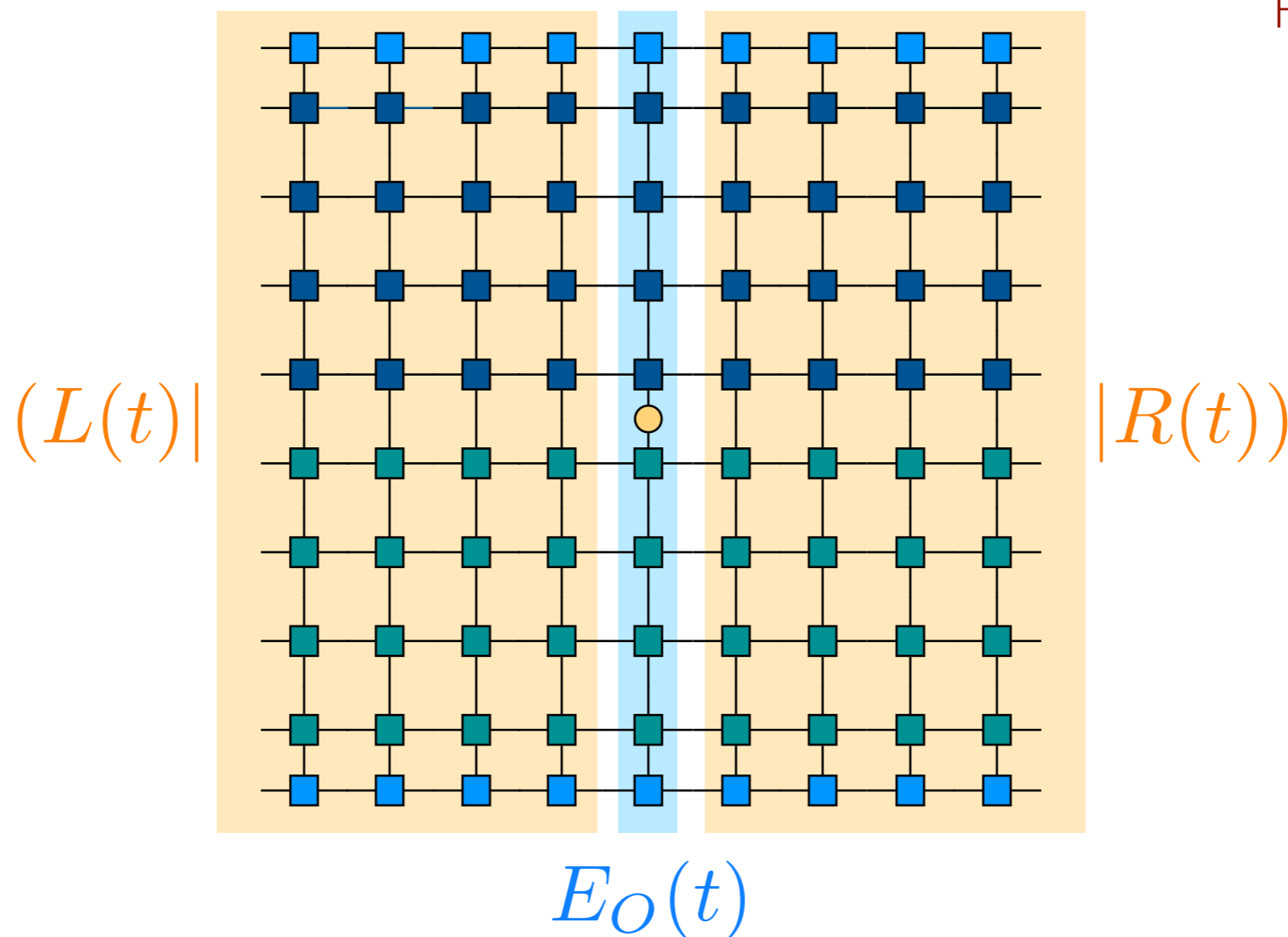
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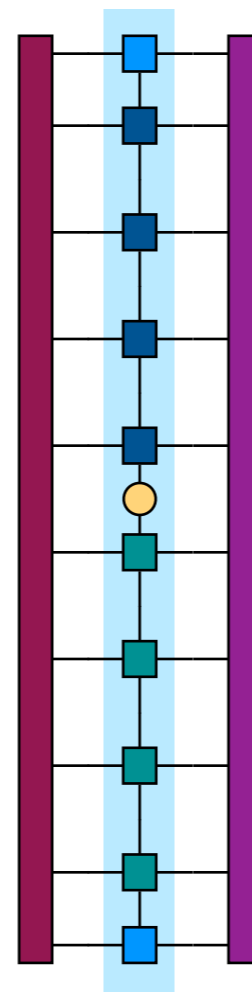
for infinite systems, transverse folding approach

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Hastings, Mahajan 2014

$|L(t)\rangle$



$|R(t)\rangle$

$E_O(t)$

transverse folding approach

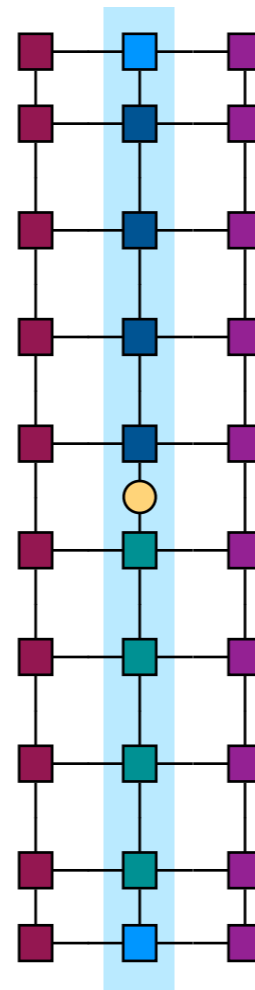
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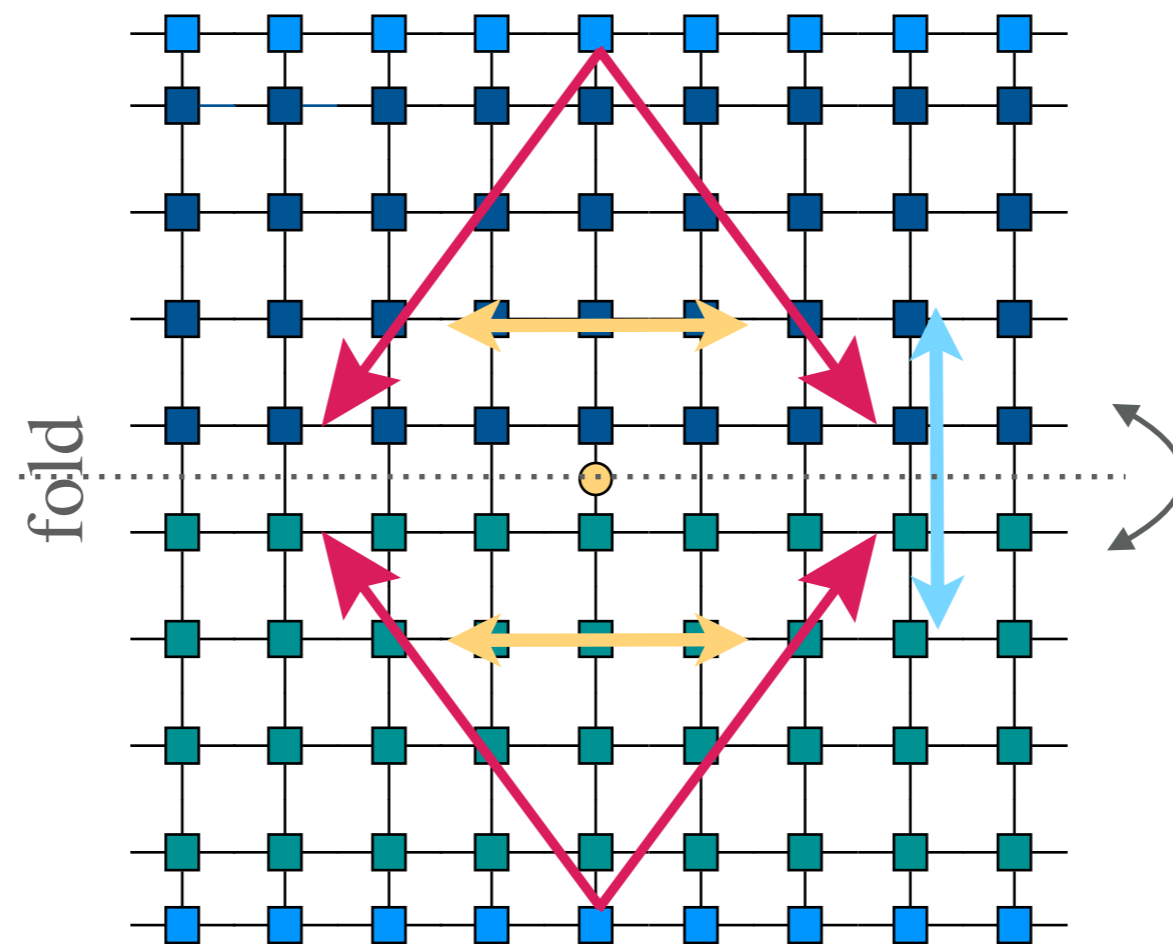
transverse folding approach

intuition: model free propagating excitations

MCB, Hastings, Verstraete, Cirac, PRL 2009

Müller-Hermes, Cirac, MCB, NJP 2012

Hastings, Mahajan 2014



transverse folding approach

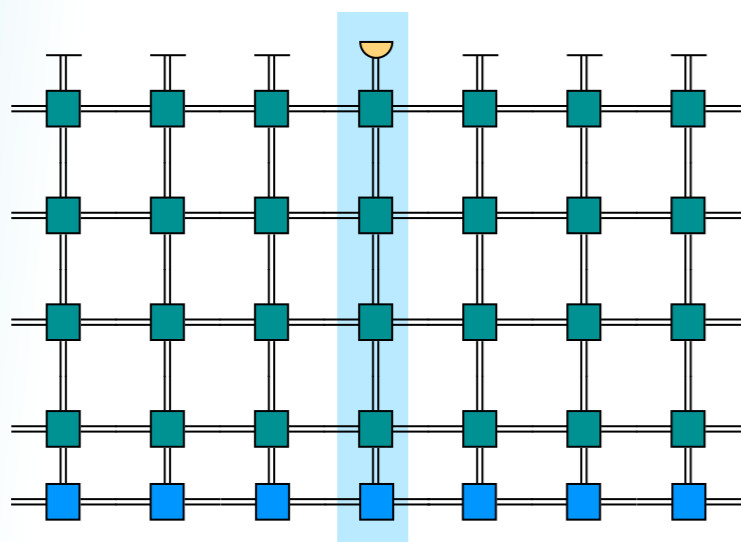
free propagating excitations

recently: influence functional
Sonner et al, Ann. Phys 2021

Lerose et al. PRX 2021

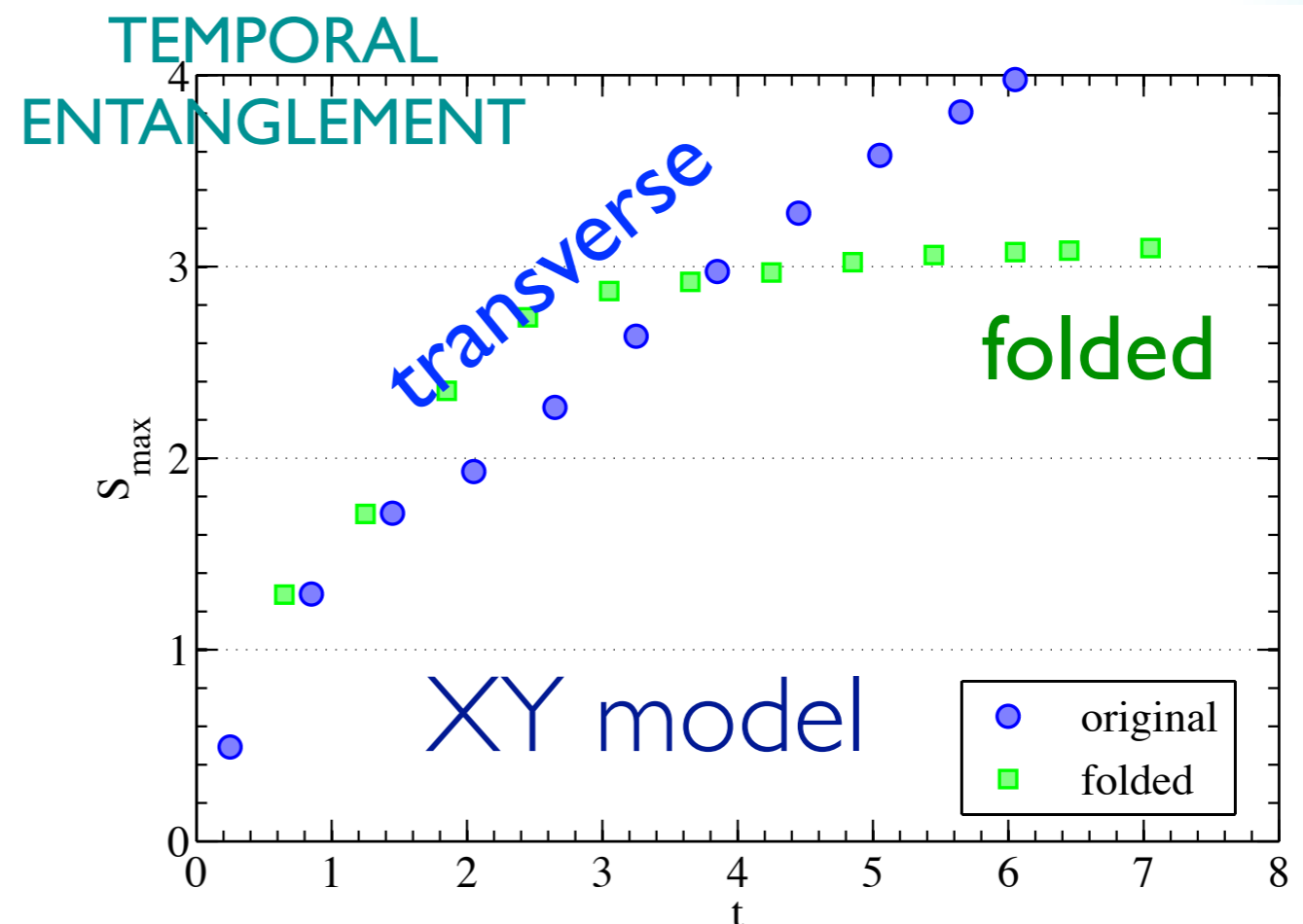
Ye, Chan, J. Chem. Phys. 2021

see also Carignano 2307.11649



$E_0(t)$

closest real case: global quench
in free fermionic models



Müller-Hermes, Cirac, MCB, NJP 2012
see also Giudice et al., PRL 128, 220401 (2022)

transverse folding + light cone = TLCC

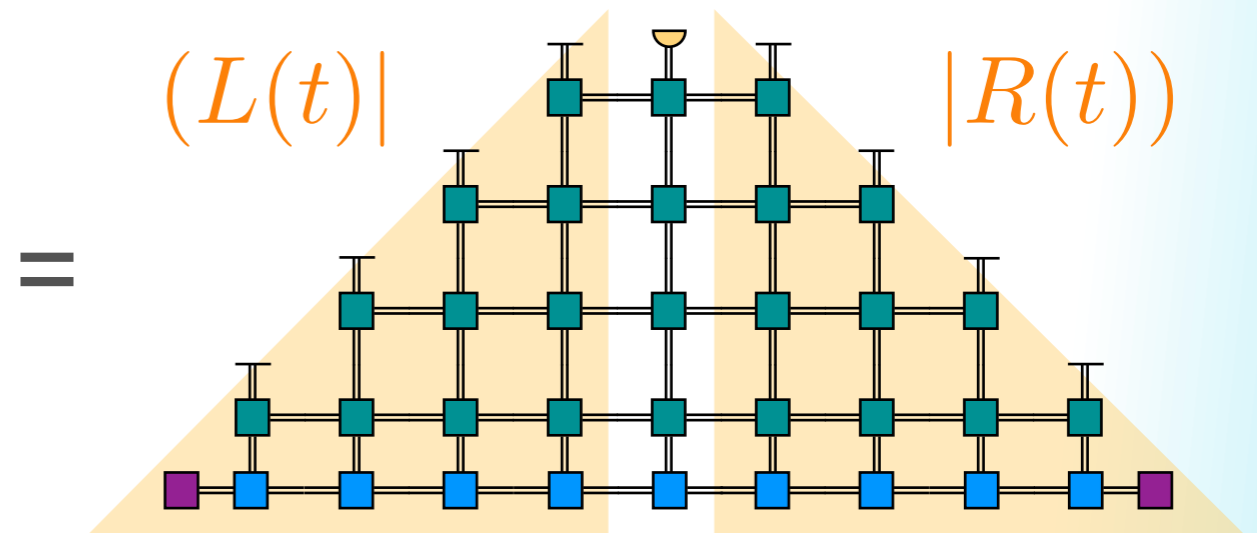
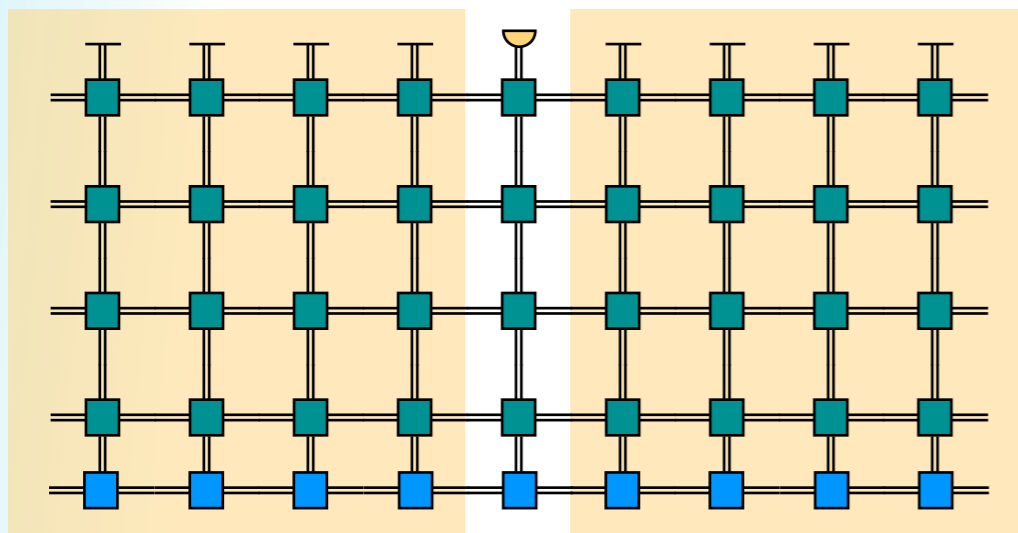
cancelling local unitaries

gain in efficiency

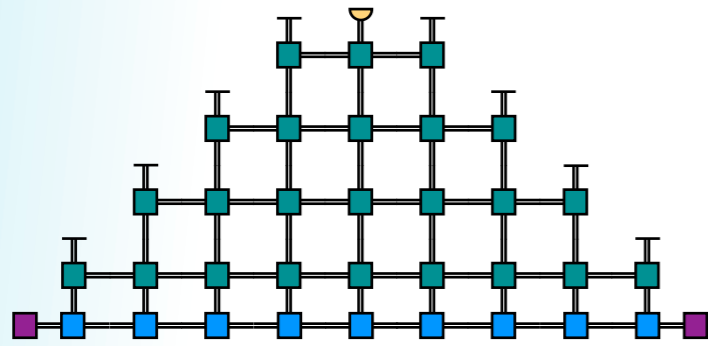
systematic increment of t

improved convergence with
Hastings' truncation

Hastings, Mahajan 2014



transverse folding + light cone = TLCC

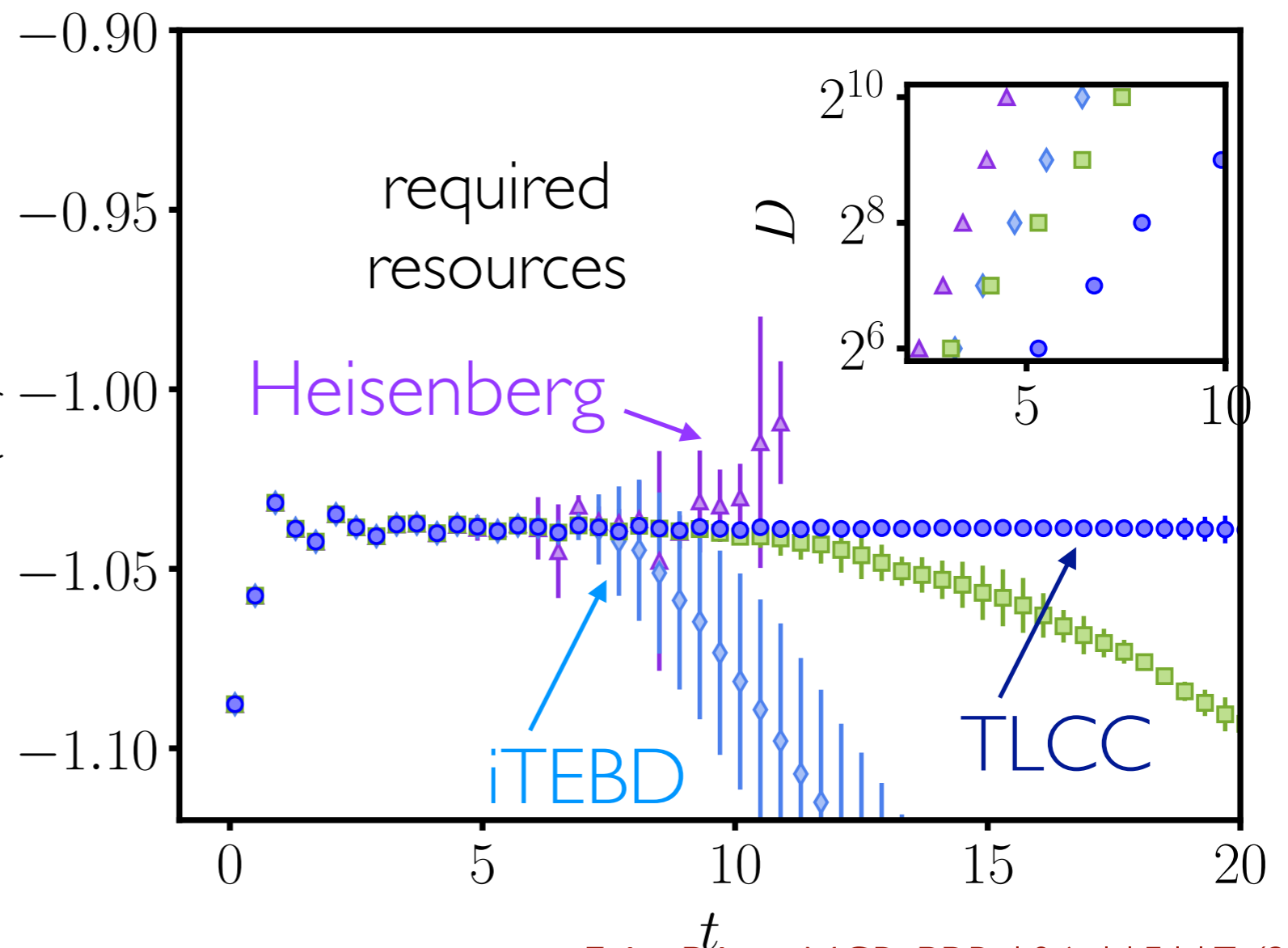


$$H_{\text{Ising}} = J \sum_{i=1}^{N-1} \sigma_z^{[i]} \sigma_z^{[i+1]} + g \sum_i^N \sigma_x^{[i]} + h \sum_i^N \sigma_z^{[i]}$$

$(J, g, h) =$

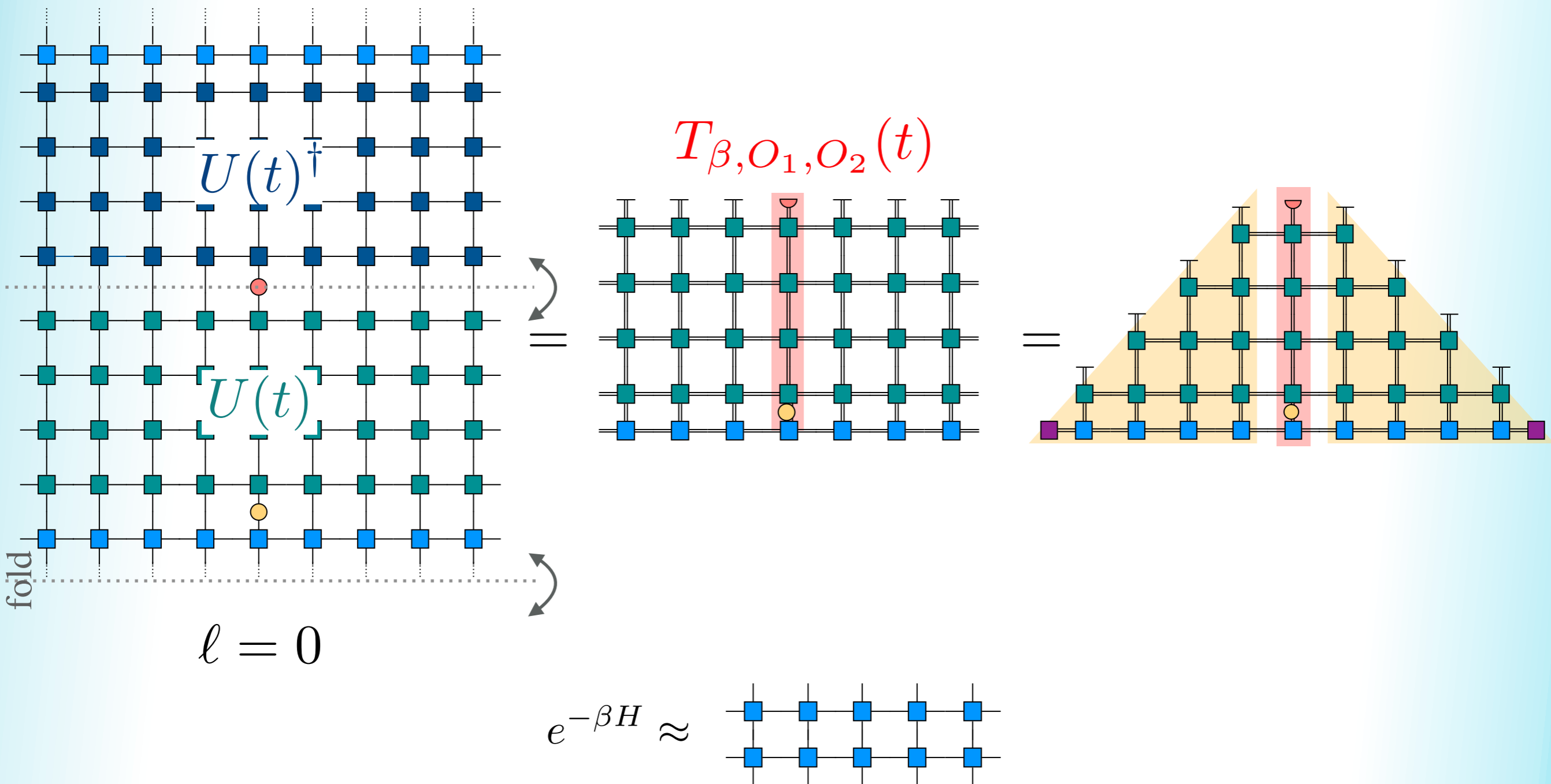
$(1, -1.05, 0.5) \langle H \rangle$

quench from $|X+\rangle$



computing response functions

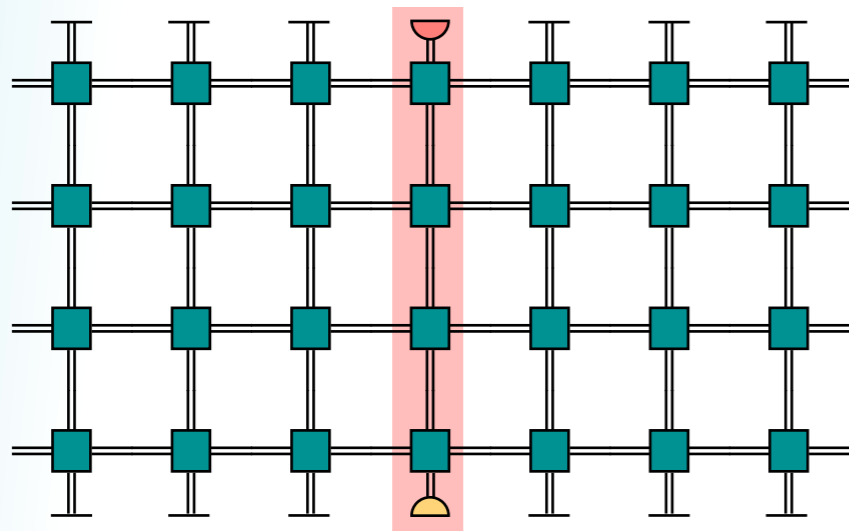
$$C_{1,2}(t, \ell, \beta) = \text{tr}(\rho_\beta O_2^{[\ell]}(t) O_1^{[0]}(0))$$



computing response functions

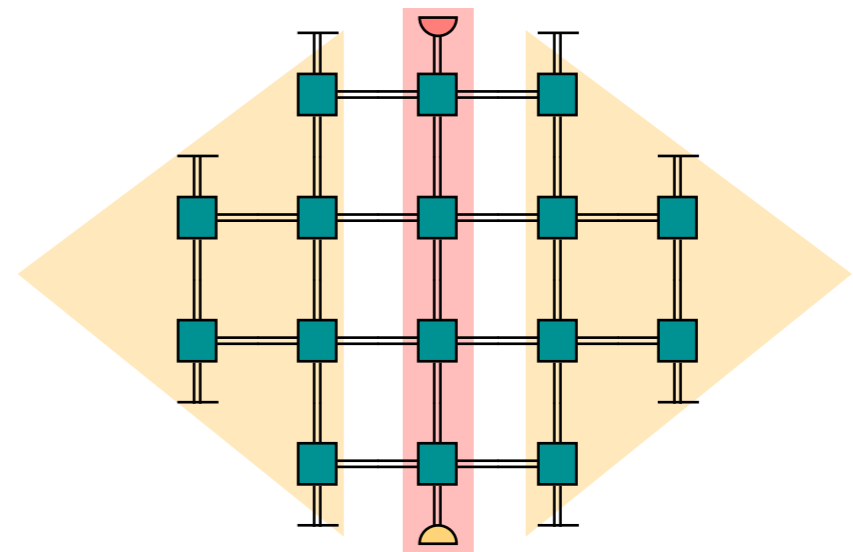
$$C_{1,2}(t, \ell, \beta) = \text{tr}(\rho_\beta O_2^{[\ell]}(t) O_1^{[0]}(0))$$

$\ell = 0$



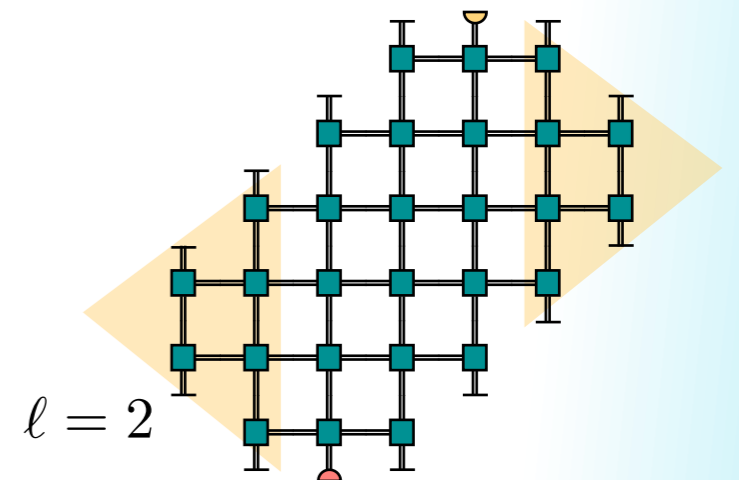
$T_{\beta=0, O_1, O_2}(t)$

=



infinite temperature

$\beta = 0$



$\ell = 2$

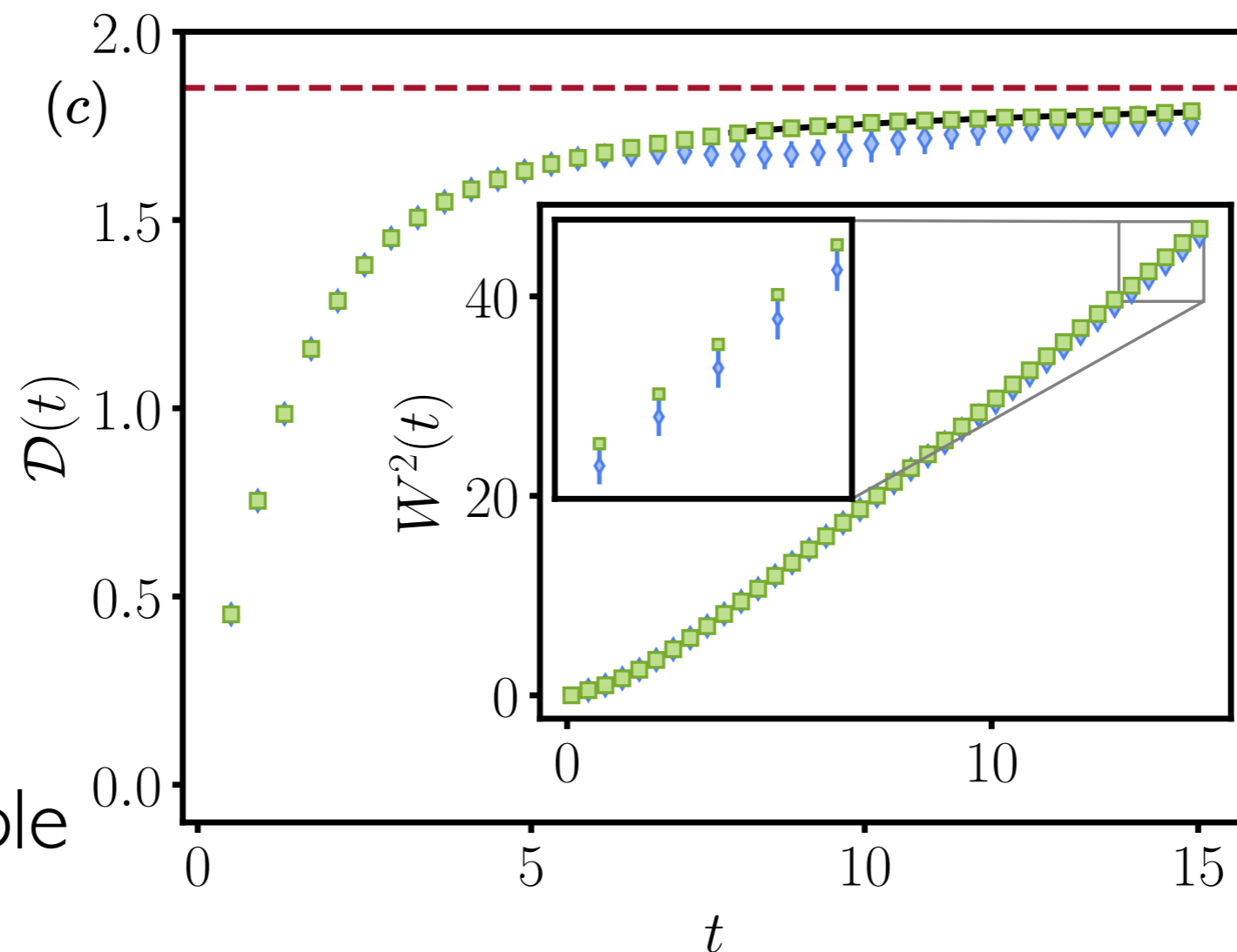
computing response functions

tilted Ising

energy density

infinite temperature

$$O_E^{[i]} := J\sigma_i^z\sigma_{i+1}^z + \frac{g}{2}(\sigma_i^x + \sigma_{i+1}^x) + \frac{h}{2}(\sigma_i^z + \sigma_{i+1}^z)$$



diffusion
constant

non-integrable
(1, -1.05, 0.5)

alternative: give up description of the full state

① light-cone TN for
non-equilibrium
evolution of local
observables

**M. Frías-Pérez, MCB,
PRB 106, 115117 (2022)**

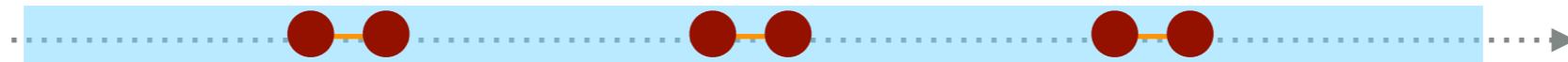
② transforming long-range
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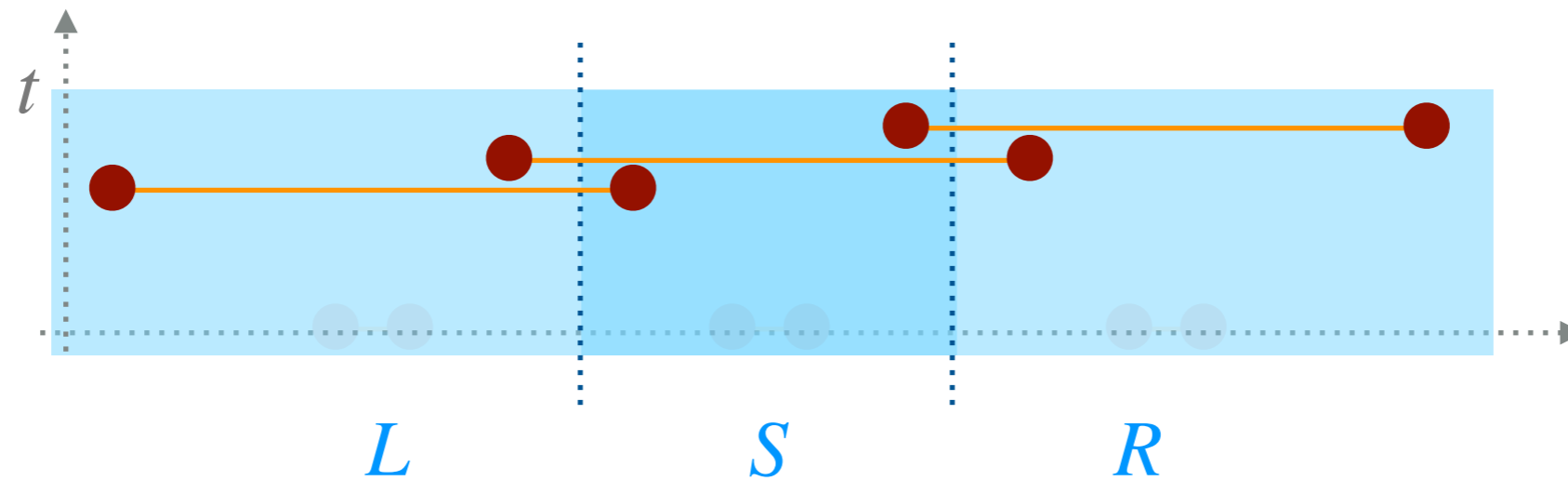
③ spectral properties of
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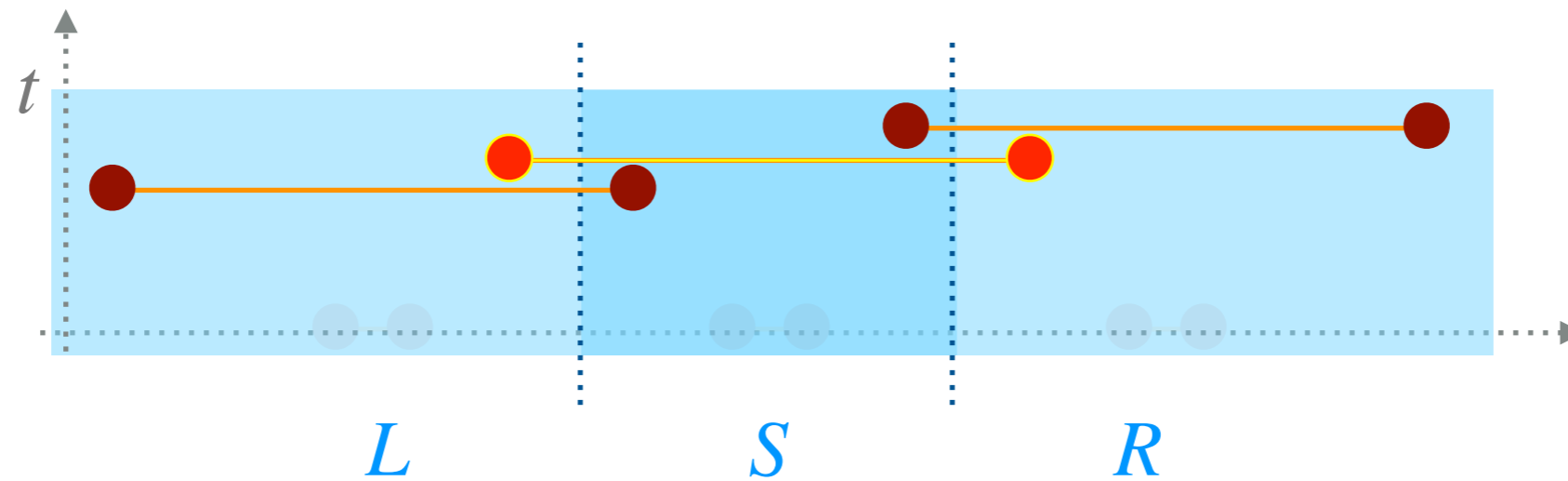
back to quasiparticle intuition



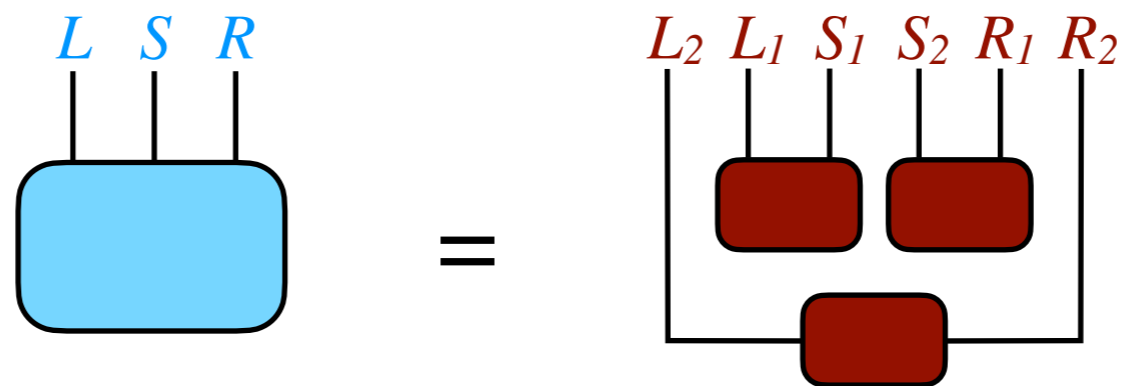
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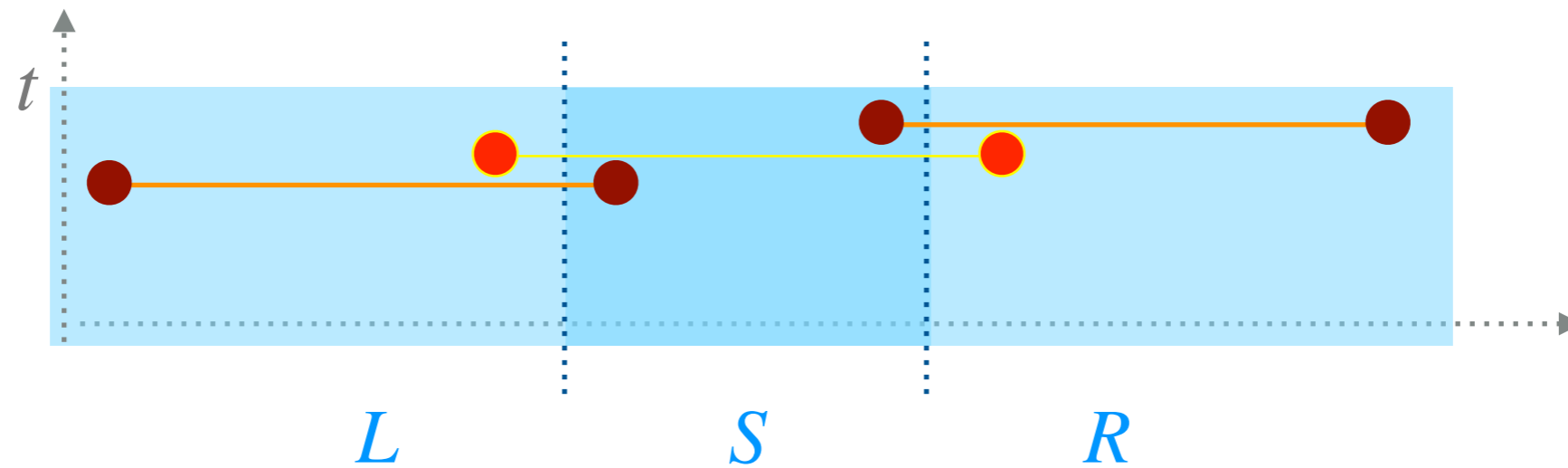
back to quasiparticle intuition



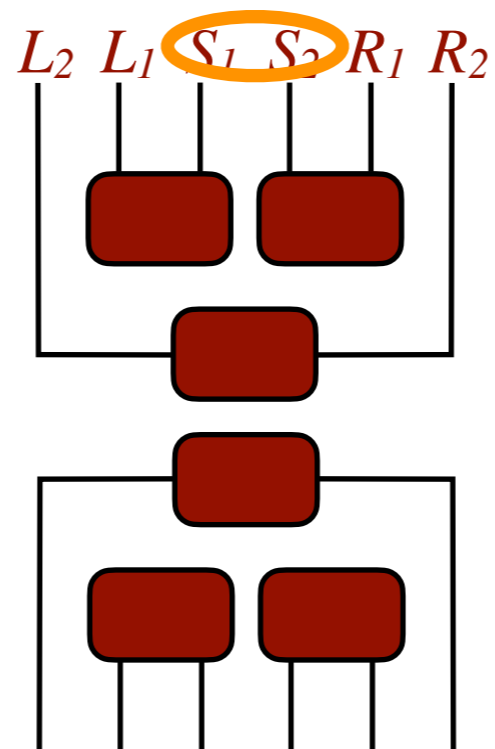
$$|\Psi_{LSR}\rangle = |\phi_{L_1 S_1}\rangle \otimes |\phi_{S_2 R_1}\rangle \otimes |\phi_{L_2 R_2}\rangle$$



back to quasiparticle intuition

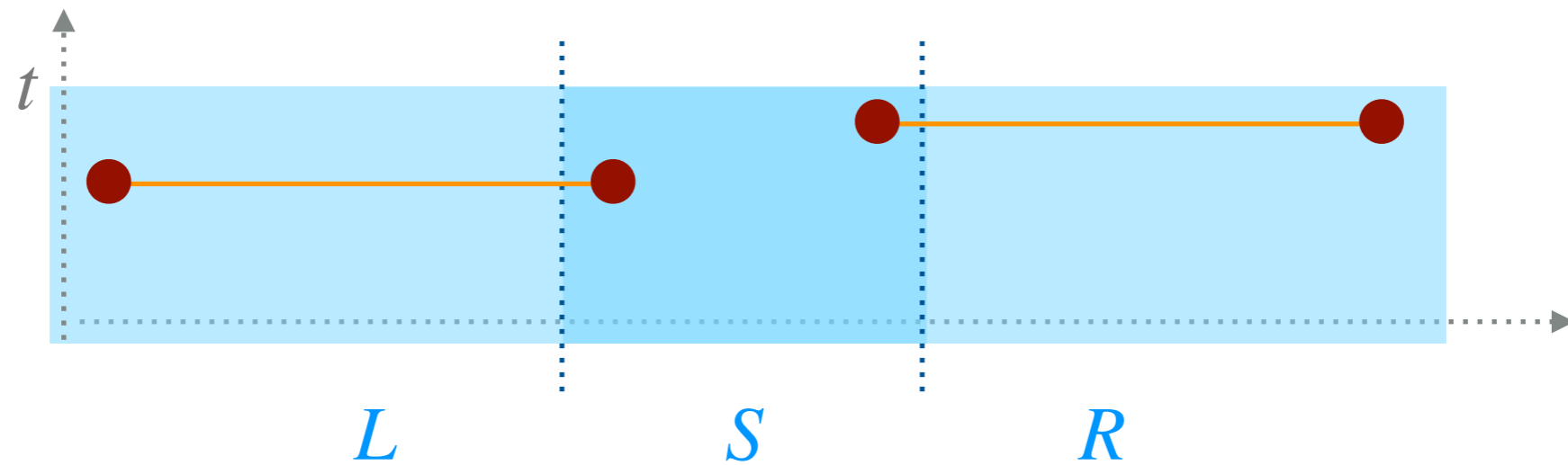


local observables on
any S block

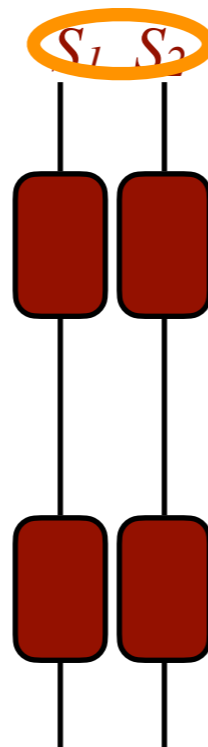


$$|\Psi_{LSR}\rangle\langle\Psi_{LSR}|$$

back to quasiparticle intuition

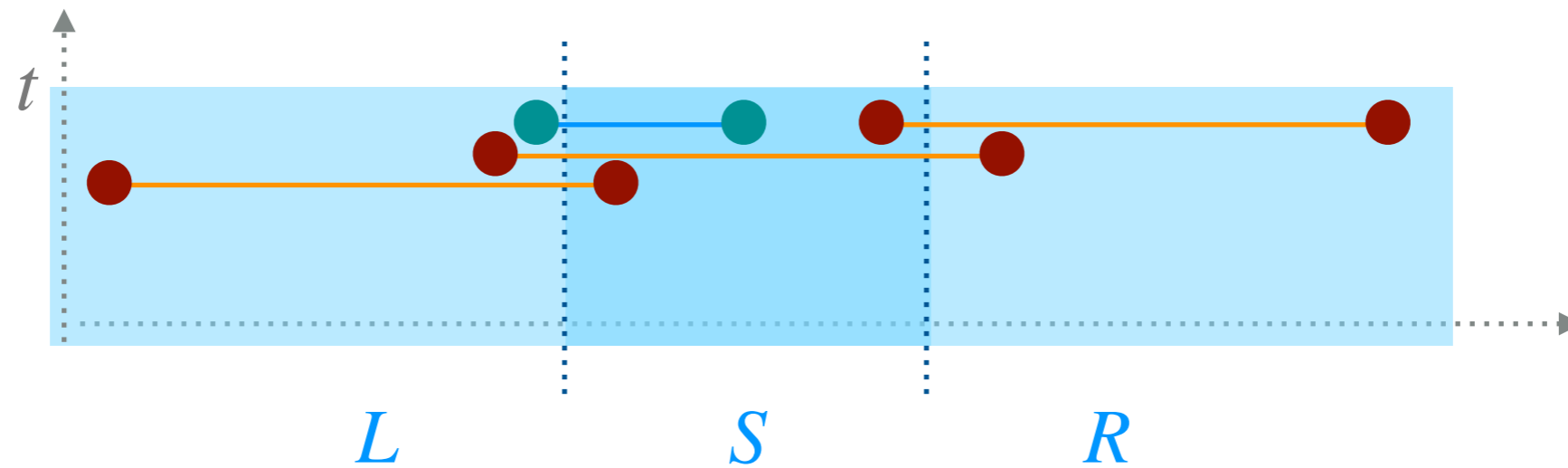


local observables on
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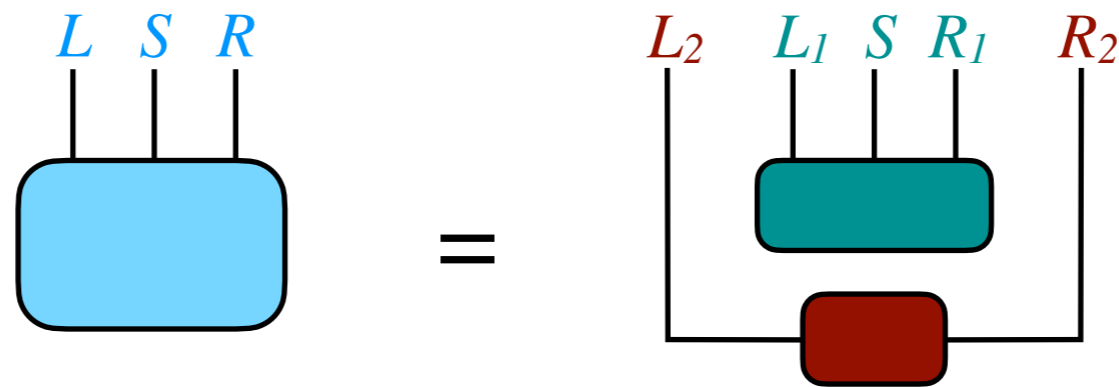


long-range
entanglement
contributes as
statistical mixture

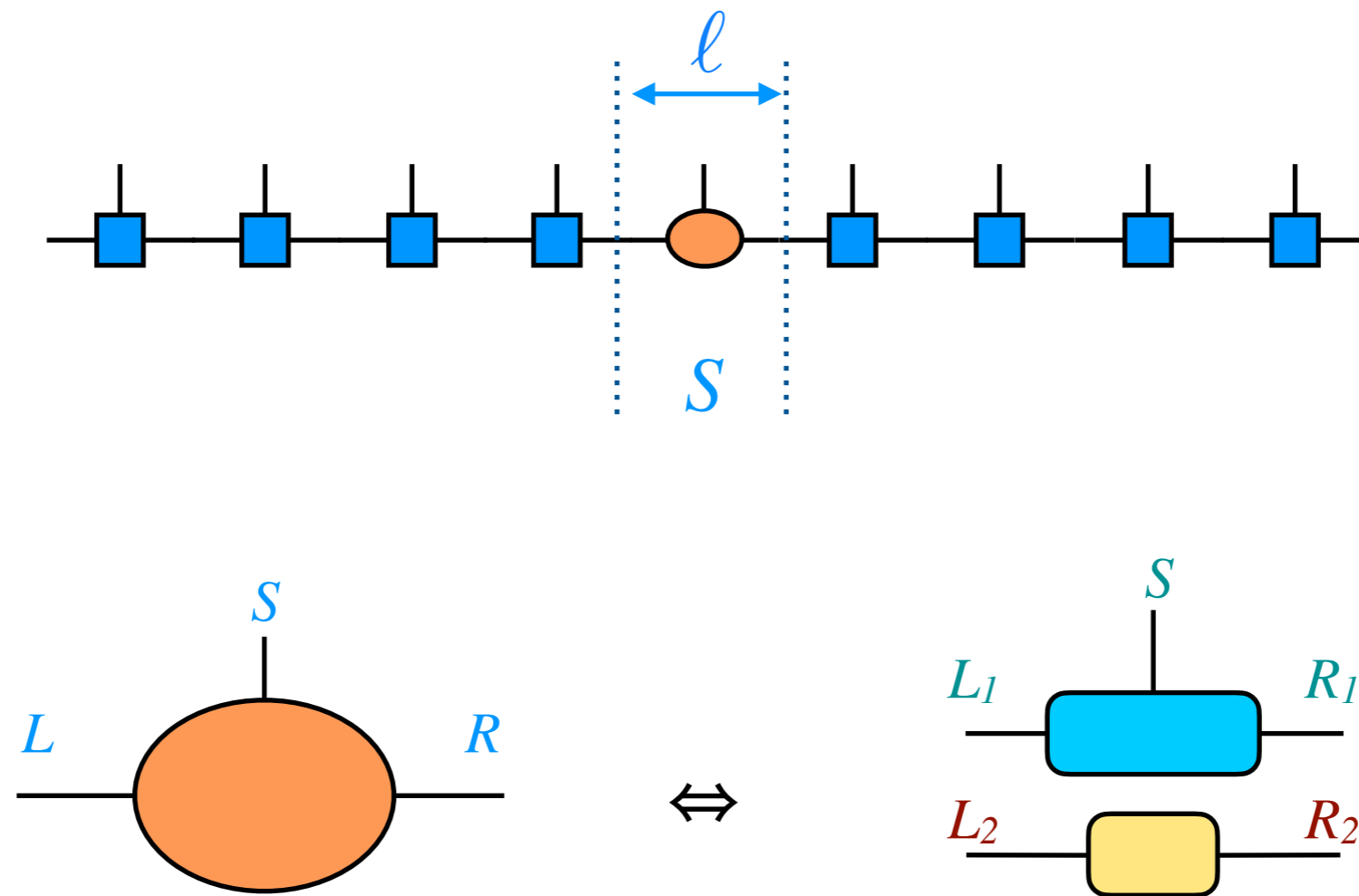
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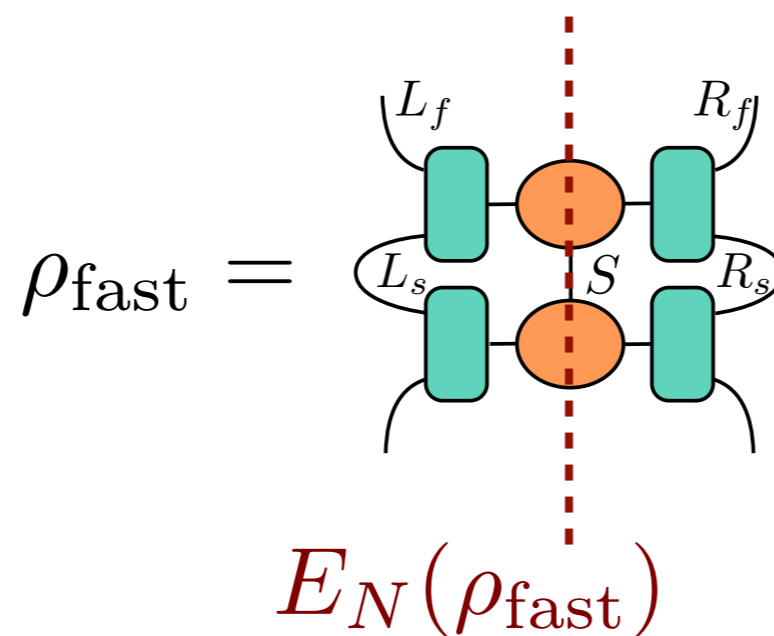
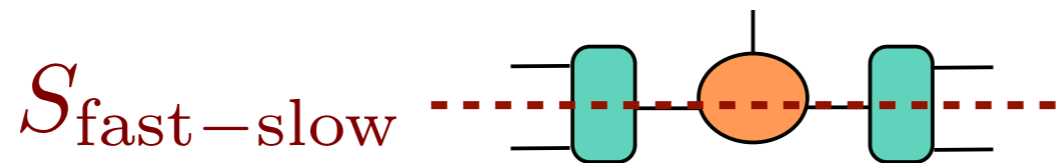
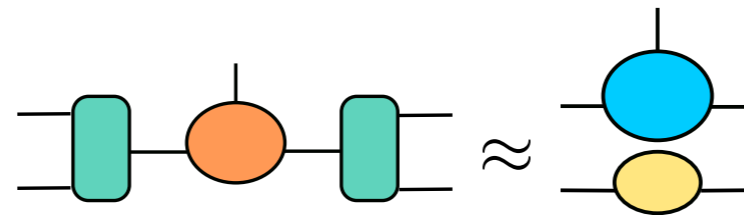
$$|\Psi_{LSR}\rangle = |\phi_{L_1SR_1}\rangle \otimes |\phi_{L_2R_2}\rangle$$



TN picture

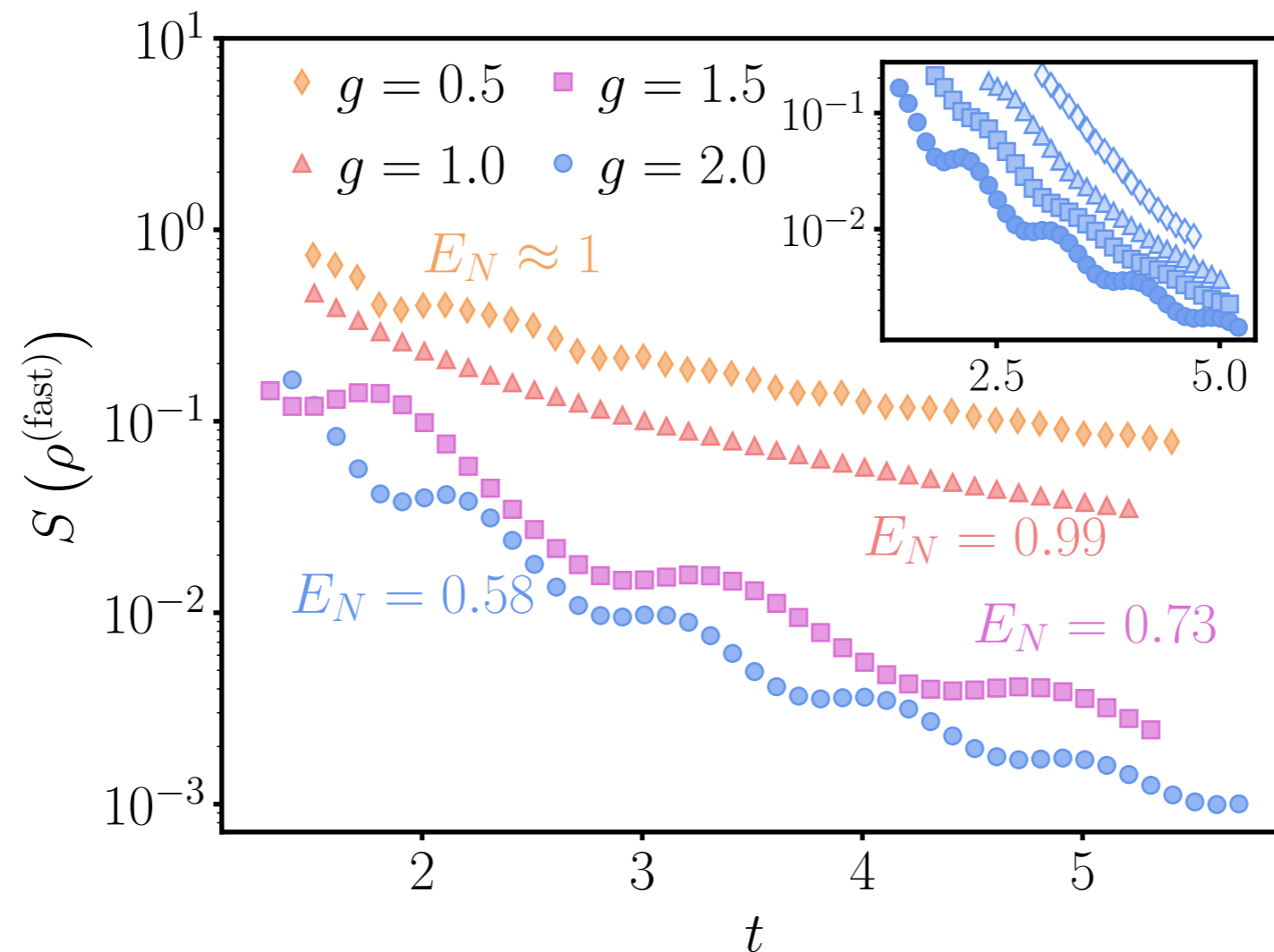
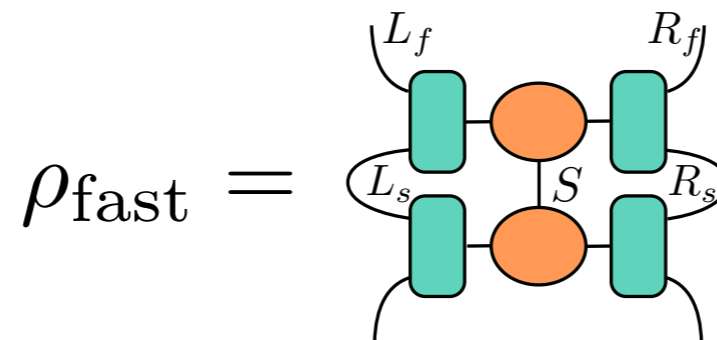


long-range entanglement in TN



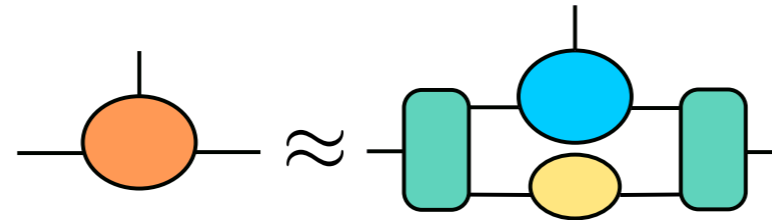
long-range entanglement in TN

$$H_{\text{Ising}} = -J \left(\sum_{i=1}^{N-1} \sigma_z^{[i]} \sigma_z^{[i+1]} + g \sum_i^N \sigma_x^{[i]} \right)$$

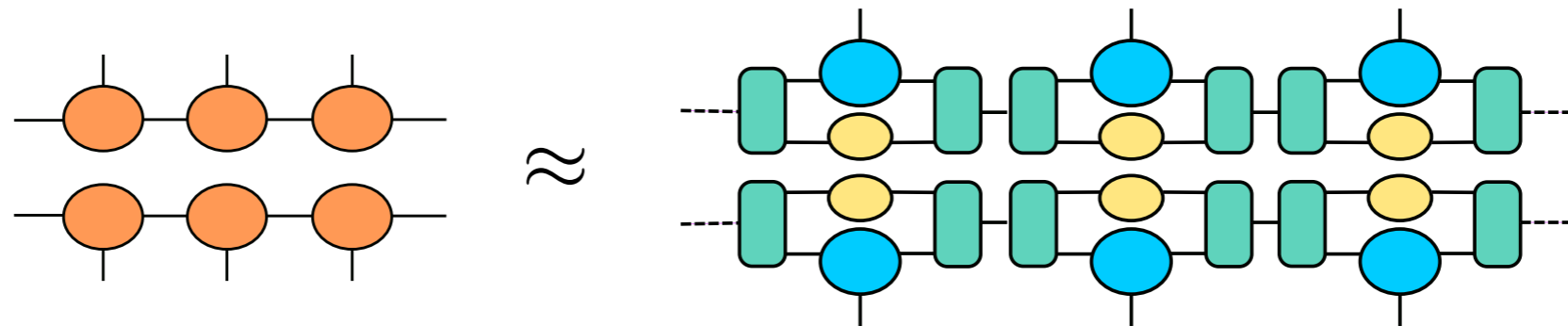


a TN algorithm

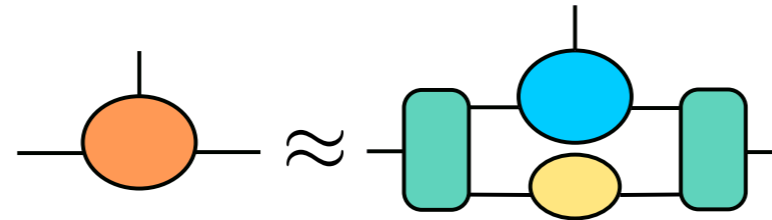
effective TN description



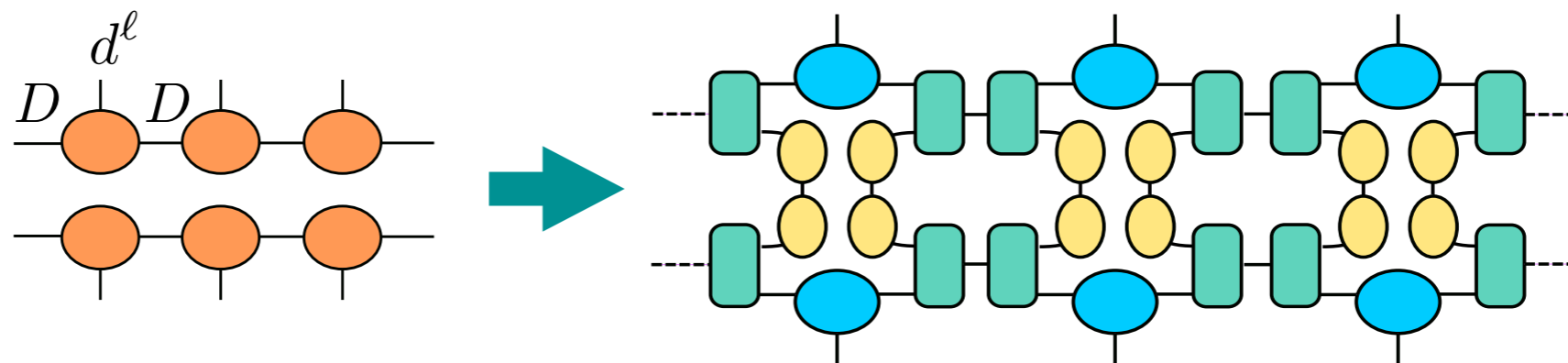
contribute as mixture to local observables in neighbouring blocks



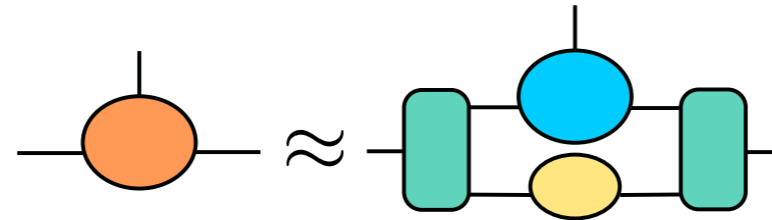
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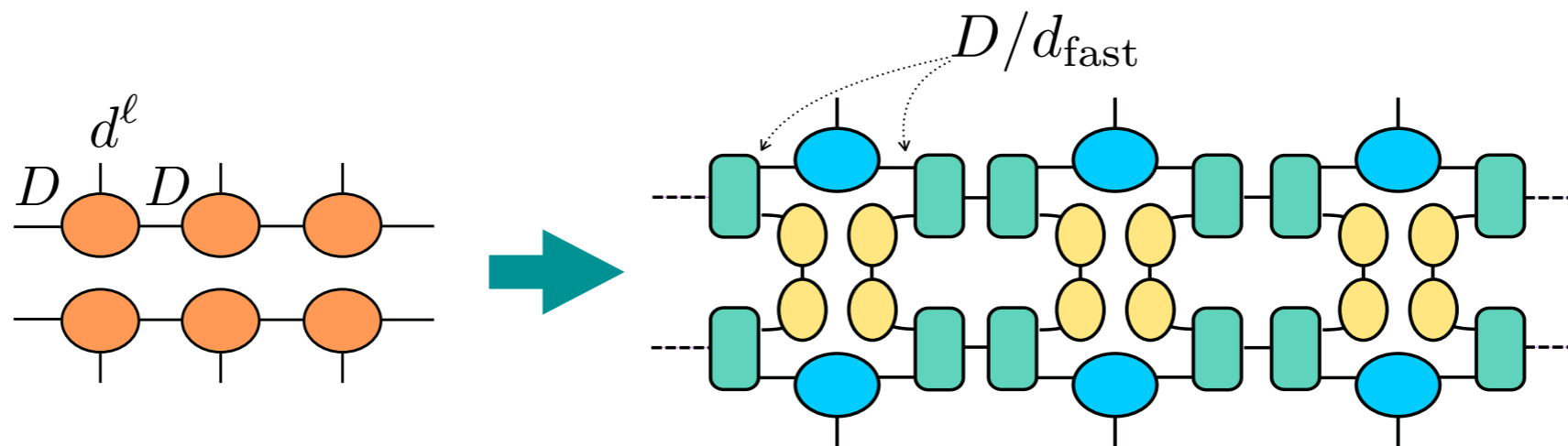
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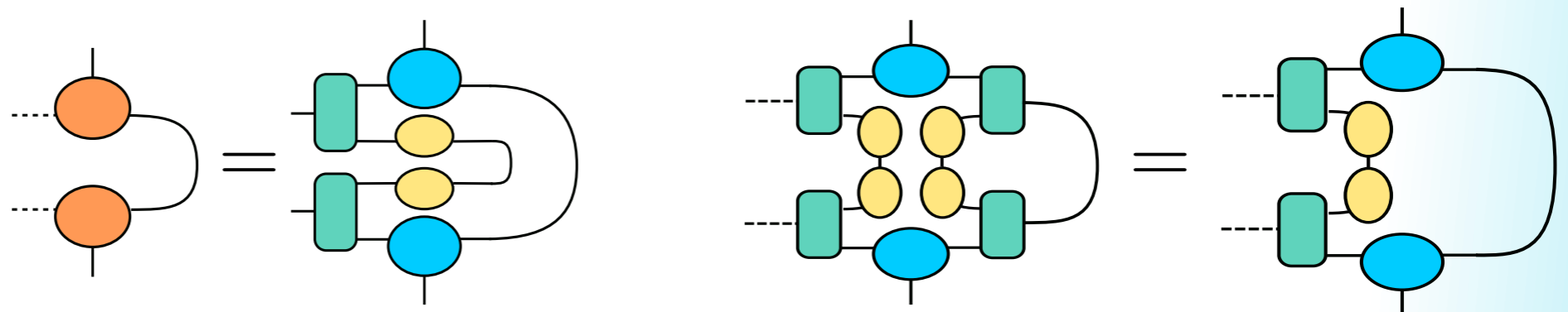
effective TN description



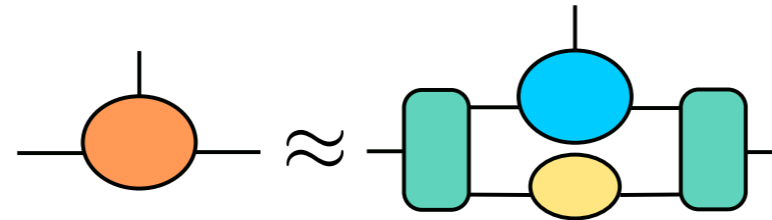
contribute as mixture to local observables in neighbouring blocks



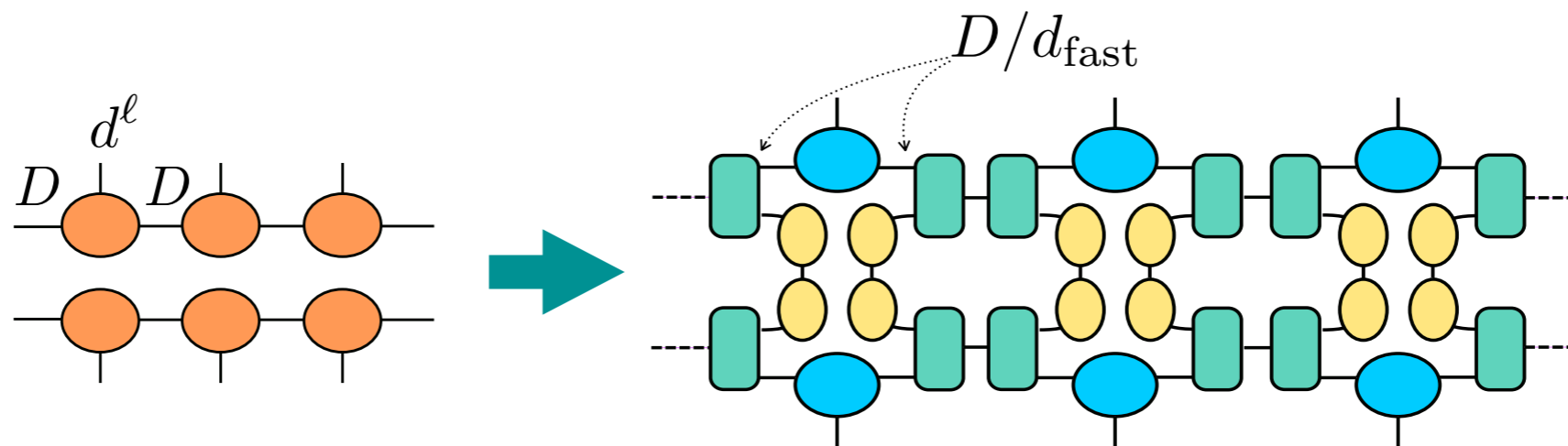
preserves reduced density matrices



effective TN description

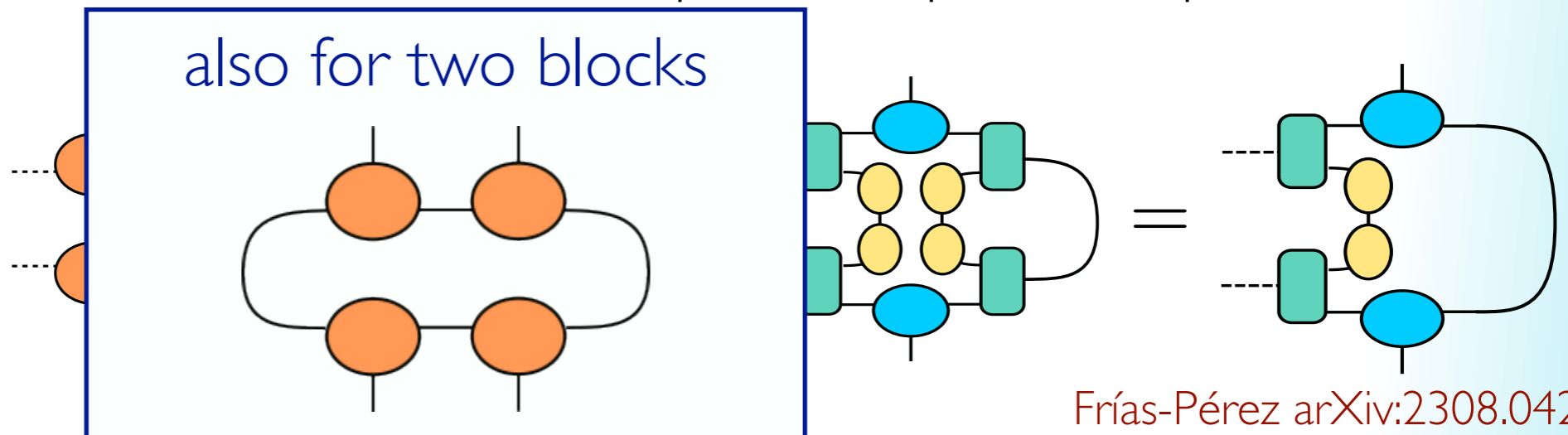


contribute as mixture to local observables in neighbouring blocks

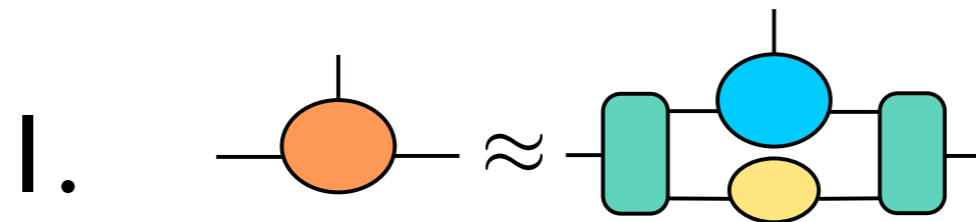


also for two blocks

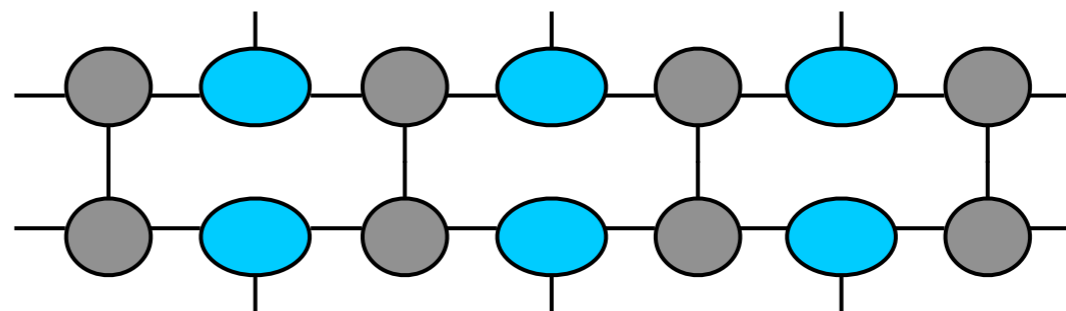
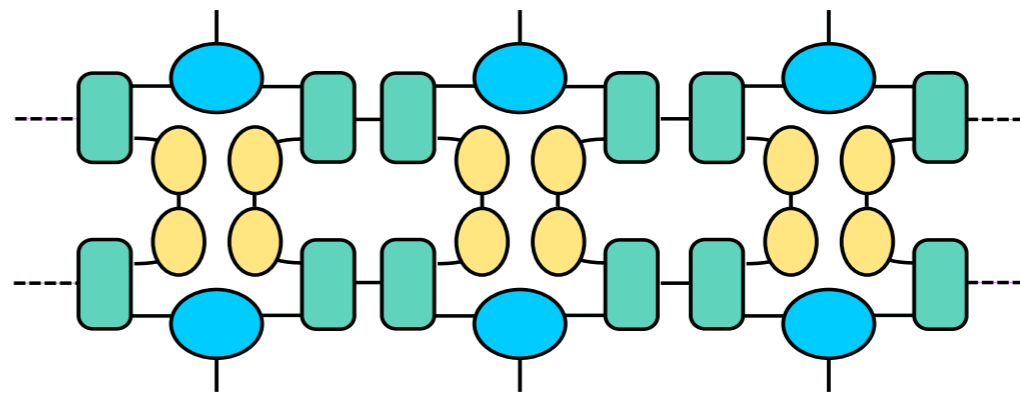
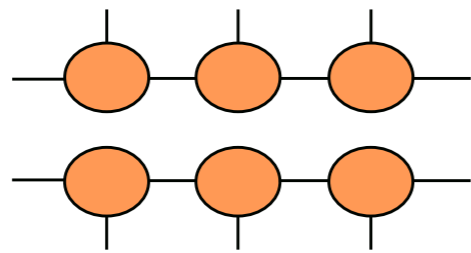
preserves reduced density matrices



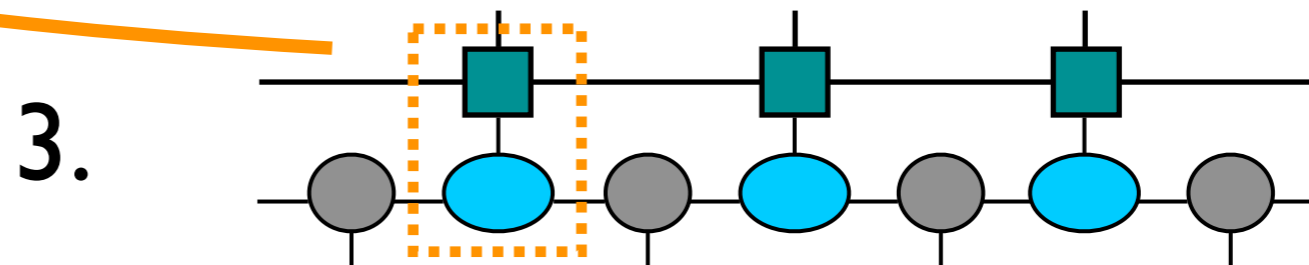
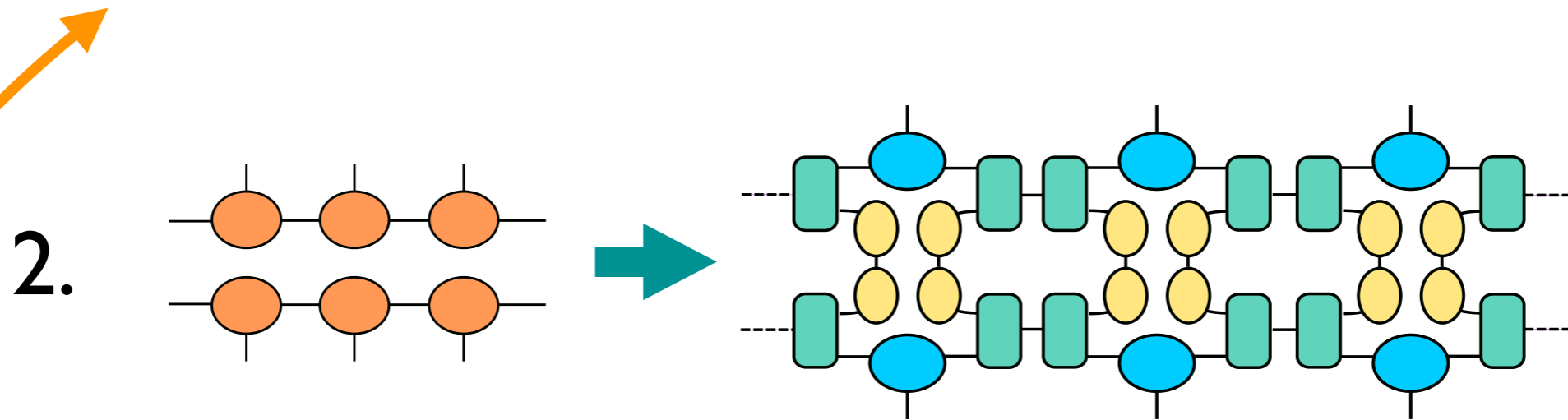
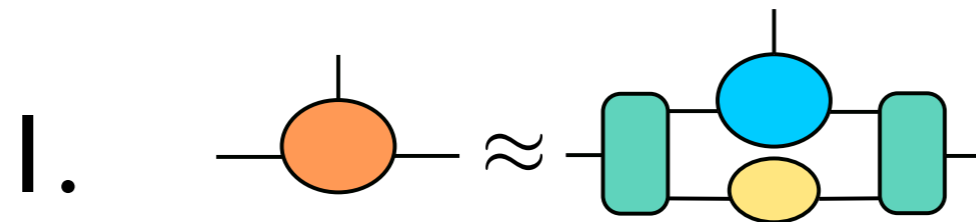
TN algorithm



2.



TN algorithm

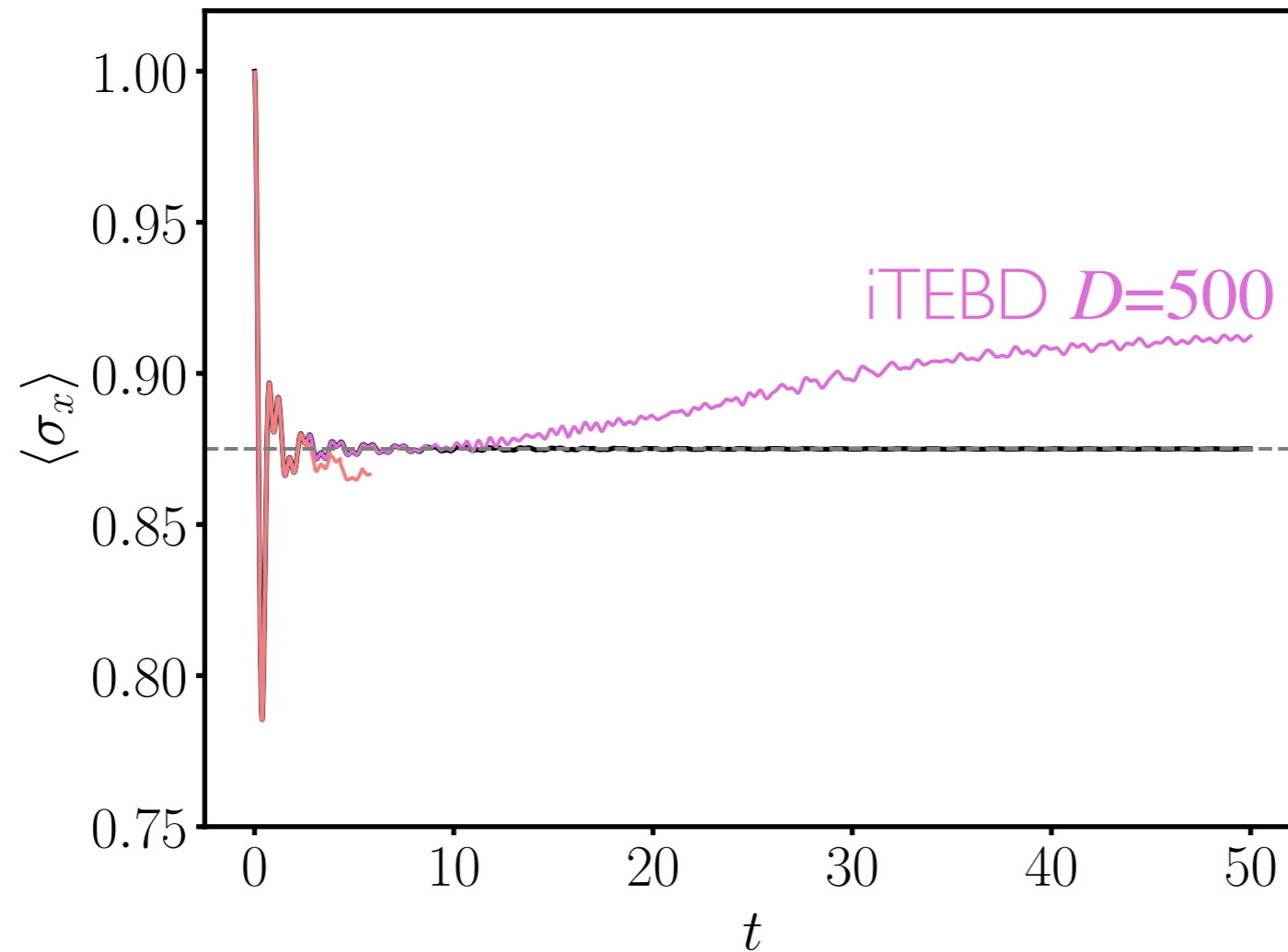


TN algorithm

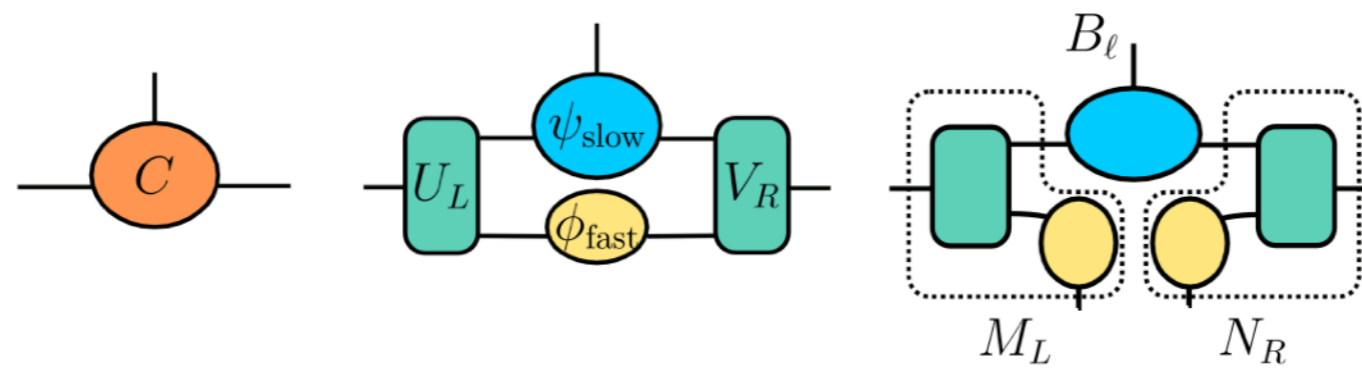
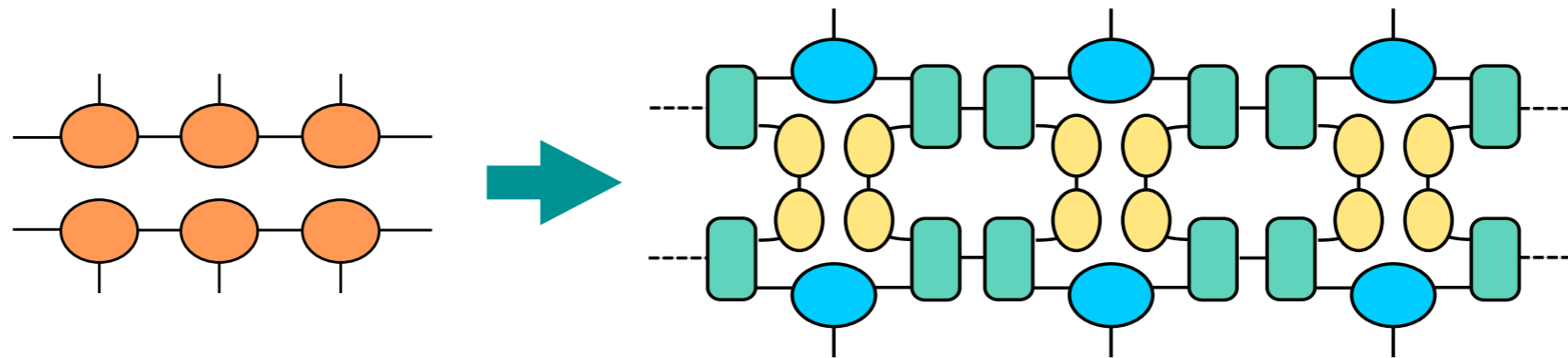
$$H_{\text{Ising}} = -J \left(\sum_{i=1}^{N-1} \sigma_z^{[i]} \sigma_z^{[i+1]} + g \sum_i^N \sigma_x^{[i]} \right)$$

$$J=1$$

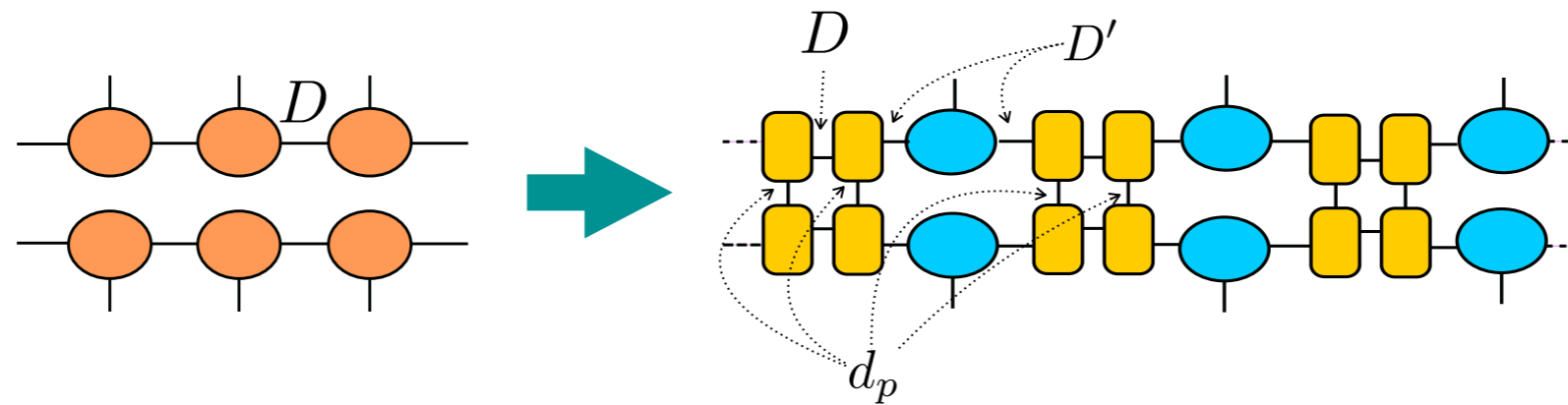
$$g=2$$



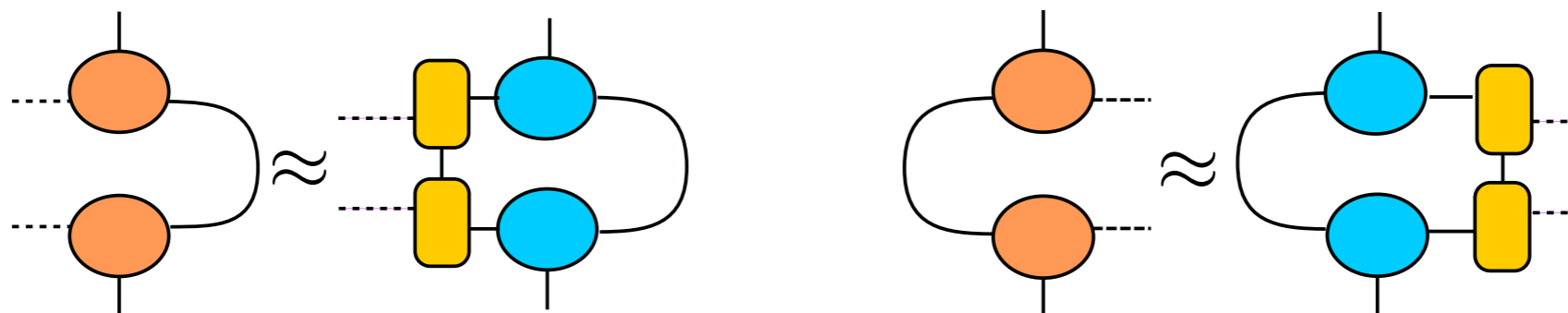
improved heuristic TN algorithm



improved heuristic TN algorithm



such that condition on rdms holds

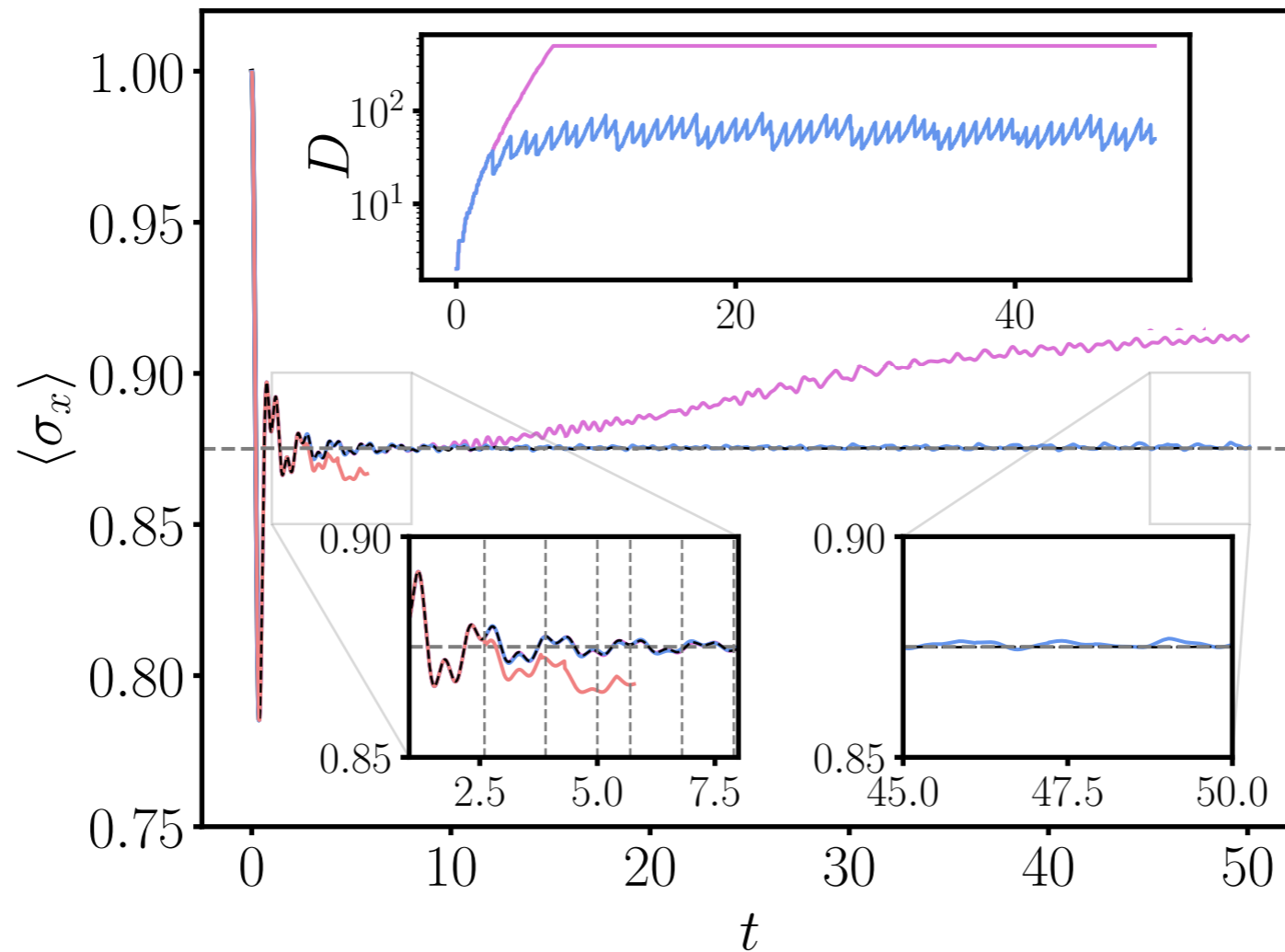


TN algorithm

$$H_{\text{Ising}} = -J \left(\sum_{i=1}^{N-1} \sigma_z^{[i]} \sigma_z^{[i+1]} + g \sum_i^N \sigma_x^{[i]} \right)$$

$$J=1$$

$$g=2$$



To conclude



energy filters & TNS can provide other
(classical and quantum) tools to get dynamical
properties

changing entanglement
perspective:
transforming long-range
entanglement into
mixture

Frías-Pérez, Tagliacozzo, MCB,
arXiv:2308.04291

light-cone TN contraction
for local observables

Frías-Pérez, MCB, PRB 106, 115117 (2022)



entanglement
barrier

spectral properties of a
QMB Hamiltonian

Yang, Iblisdir, Cirac, MCB, PRL 124, 100602 (2020)

Lu, PRX Quantum 2, 020321 (2021)

Yang, Cirac, MCB, PRB 106, 024307 (2022)

